

# 17 O relaxation times in the rat brain at 16.4 tesla

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# $^{17}\text{O}$ Relaxation Times in the Rat Brain at 16.4 Tesla

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**Purpose:** Measurement of the cerebral metabolic rate of oxygen (CMRO<sub>2</sub>) by means of direct imaging of the  $^{17}\text{O}$  signal can be a valuable tool in neuroscientific research. However, knowledge of the longitudinal and transverse relaxation times of different brain tissue types is required, which is difficult to obtain because of the low sensitivity of natural abundance  $\text{H}_2^{17}\text{O}$  measurements.

**Methods:** Using the improved sensitivity at a field strength of 16.4 Tesla, relaxation time measurements in the rat brain were performed in vivo and postmortem with relatively high spatial resolutions, using a chemical shift imaging sequence.

**Results:** In vivo relaxation times of rat brain were found to be  $T_1 = 6.84 \pm 0.67$  ms and  $T_2^* = 1.77 \pm 0.04$  ms. Postmortem  $\text{H}_2^{17}\text{O}$  relaxometry at enriched concentrations after inhalation of  $^{17}\text{O}_2$  showed similar  $T_2^*$  values for gray matter ( $1.87 \pm 0.04$  ms) and white matter, significantly longer than muscle ( $1.27 \pm 0.05$  ms) and shorter than cerebrospinal fluid ( $2.30 \pm 0.16$  ms).

**Conclusion:** Relaxation times of brain  $\text{H}_2^{17}\text{O}$  were measured for the first time in vivo in different types of tissues with high spatial resolution. Because the relaxation times of  $\text{H}_2^{17}\text{O}$  are expected to be independent of field strength, our results should help in optimizing the acquisition parameters for experiments also at other MRI field strengths. **Magn Reson Med 000:000–000, 2015. © 2015 Wiley Periodicals, Inc.**

**Key words:** relaxation times; brain water distribution; natural abundance oxygen-17;  $\text{H}_2^{17}\text{O}$ ; quadrupolar relaxation; X-Nuclei; quantification; spectroscopic imaging; MRSI; ultra-high field

emission tomography (6)], MRI during catabolism of enriched  $^{17}\text{O}_2$  to  $\text{H}_2^{17}\text{O}$  can become a valuable alternative. However, the NMR detection of  $^{17}\text{O}$  is challenging due to its low gyromagnetic ratio (7), short transverse relaxation times (8) and low natural abundance (9). The sensitivity of  $^{17}\text{O}$  NMR has been shown to increase almost quadratically with field strength, supported by the field-independent relaxation times of unbound  $\text{H}_2^{17}\text{O}$  (8,10,11). Further improvements can be realized by optimizing the acquisition parameters based on the relaxation times (12,13).

The relaxation mechanisms of the  $^{17}\text{O}$  nucleus are significantly different from those of protons due to the spin of  $+5/2$  (7), generally being more intramolecular as a result of quadrupolar interaction. Water as the exclusive source of the in vivo NMR signal of  $^{17}\text{O}$  as triplet (14–16) exhibits one peak in bulk water due to fast hydrogen exchange and isotropic molecular motion (17). At physiological conditions around neutral pH, the relaxation of  $\text{H}_2^{17}\text{O}$  is strongly influenced by chemical exchange and by temperature (15). Because field-dependent relaxation has been reported for bound water (18), a regional determination of representative relaxation times in complex in vivo tissue as a mixture of bound and free water in the brain is important for applying CMRO<sub>2</sub> imaging through  $\text{H}_2^{17}\text{O}$  NMR detection.

In this study, we determined the regional longitudinal  $T_1$  and apparent transverse  $T_2^*$  relaxation times of  $\text{H}_2^{17}\text{O}$  at natural abundance and in enriched concentrations in the rat head at 16.4T. Due to the high magnetic field and optimized measurement parameters, local differences in the relaxation times among different tissue types of rats have been observed for the first time in vivo at relatively high spatial resolutions. Additional phantom measurements were performed to assess the changes of the relaxation times with the temperature, which is critical due to the influence of viscosity and chemical exchange (19).

Previous relaxometric studies of  $\text{H}_2^{17}\text{O}$  in rat brain have not differentiated between different tissue types in vivo. Finally, because pH, temperature, and concentrations of salts affect the chemical exchange and molecular motion, the linewidth of  $\text{H}_2^{17}\text{O}$  is naturally autoregulated in the brain within a narrow range. Therefore, the advantages of ultra-high field MRI make it possible to obtain the necessary signal sampling efficiency at sufficient spatial resolutions and the observed field independence of the  $\text{H}_2^{17}\text{O}$  relaxation times ensures that these results can easily be applied to studies at other field strengths.

## METHODS

### Instrumentation

Experiments were performed on a Bruker BioSpec Avance III system (Bruker Biospin MRI GmbH, Ettlingen,

## INTRODUCTION

In vivo imaging of oxygen metabolism, a key biological parameter for characterizing the vital status of tissue, can have a large impact in medicine and neuroscience. While most established techniques used to measure the cerebral metabolic rate of oxygen (CMRO<sub>2</sub>) have the limitations of being either indirect (e.g., calibrated BOLD) (1,2) or invasive [arteriovenous difference measurements (3–5) and positron

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Germany) based on a 16.4 Tesla (T) magnet (Magnex Scientific, Abingdon, UK) with a 26-cm bore, equipped with gradients with an inner diameter of 12 cm and a maximum strength of 1 T/m within a rise time of 212  $\mu$ s (Resonance Research Inc., Billerica, MA). Custom-built  $^{17}\text{O}$  quadrature surface coils (diameter 1.5 cm; sensitive detection depth  $\sim$  1 cm) (20) were tuned and matched for every sample individually to the  $^{17}\text{O}$  Larmor-frequency of 94.6 MHz. The signals from 0° and 90° channels were combined by a quadrature hybrid and then amplified and digitized by a single-channel X-band receiver. Mutual coupling between the two channels, assessed by S21 measurements, was better than -20 dB. For proton imaging, a separate  $^1\text{H}$  transmit and receive butterfly coil was attached on top of the  $^{17}\text{O}$  coils.

### Animal Preparation

All procedures and experiments were approved by the local authorities and carried out by FELASA qualified personnel in compliance with the guidelines of the European Community (EUVD 86/609/EEC) for the care and use of laboratory animals. A total of 14 male Wistar rats (Charles River Laboratories, Sulzfeld, Germany) were used in this study. Animals were kept at a 12/12 h day/night cycle with free access to nutrition and water.

### In Vivo Measurements

Ten male Wistar rats were used for in vivo relaxometry measurements ( $T_1$ ;  $n = 5$ ; mean body weight  $513 \pm 71$  g;  $T_2^*$ ;  $n = 5$ ; mean body weight  $529 \pm 30$  g). Anesthesia was induced in an induction chamber by 3–4% isoflurane (Forene®, Abbott, Chicago, IL). During the preparation of the animal, anesthesia was maintained by spontaneous inhalation of room air with 2% isoflurane through an inhalation mask. A 22-gauge catheter (Venisystems Abboath®, Abbott Laboratories, UK) was inserted into a lateral tail vein for intravenous infusion of medetomidine anesthesia. The rat head was fixated in a custom-built stereotaxic frame with the center of maximum sensitivity of the  $^{17}\text{O}$ -coil in Z-direction close to the bregma of the rat skull. Lidocain spray (Xylocain®, AstraZeneca GmbH, Wedel, Germany) was applied topically before installation of the ear-bars. A subcutaneous bolus of Flunixin-Meglumin (5 mg/kg; Finadyne®, Essex Pharma GmbH, München, Germany) was given as analgesic for the intravenous catheter and fixation. The transition to continuous medetomidine anesthesia (Domitor®, Pfizer Inc., New York, NY) was initiated by intravenous infusion at a rate of 0.1 mg/kg/h (dilution: 0.1 mg in 1-mL saline solution with 16 IU/mL of heparin [heparinatrium 25000®, ratiopharm GmbH, Ulm, Germany]) with successive discontinuation of isoflurane over 15–20 min (21). For the maintenance of anesthetic efficacy during long experiments (22), the medetomidine infusion rate was increased every 75 min by 50  $\mu$ g/kg/h.

Breathing rate and exhaled respiration gases (Capstar-100 End-tidal  $\text{CO}_2$  Analyzer, CWE Inc., Ardmore, PA) were monitored and recorded (PowerLab & LabChart 5; ADInstruments, Australia). Rectally measured body temperature was maintained at  $36.8 \pm 0.5^\circ\text{C}$  by an electric heating pad connected through a radiofrequency (RF) filter. After the

experiments, anesthesia was revoked using atipamezolhydrochlorid (Antisedan®, Pfizer Inc., New York, NY).

### Postmortem Experiments

Four male Wistar rats (mean body weight  $380 \pm 131$  g) were initially anesthetized with 3–4% and kept at 2% isoflurane during preparations. The animals were orally intubated for mechanical ventilation and intravenously catheterized for infusion of alpha-chloralose at 50 mg/kg/h. Rats were positioned inside the magnet and were subjected to a total approximately 90 min inhalation of 30%  $\text{O}_2$  with a  $^{17}\text{O}_2$ -enrichment of 70% (Nukem GmbH, Germany) in  $\text{N}_2\text{O}$  (one rat in  $\text{N}_2$ ) as part of a different study. Two hours after inhalation, the animals were killed with >5% isoflurane without changing their position within the magnet. Acquisitions started 10 min after exitus with the rat body allowed to equilibrate to ambient temperature in the magnet bore ( $21^\circ\text{C}$  at isocenter) (23).

### MRI Experiments

#### Adjustments

The RF-power for a nominal 90° excitation at the  $^{17}\text{O}$  frequency was calibrated globally for each individual sample by maximizing the signal of a “pulse-and-acquire” sequence, using a 50  $\mu$ s on-resonance block-pulse. For anatomical imaging and shimming, a proton surface coil was mounted on top of the  $^{17}\text{O}$  coil assembly, while in some of the in vivo measurements a  $^1\text{H}$  resonator surrounding the setup was used instead. First and second order  $B_0$  shims were adjusted using FASTMAP (24) with a 1-ms Gaussian pulse. Proton images were acquired at high in-plane resolutions with equal slice thickness as the corresponding  $\text{H}_2^{17}\text{O}$  images.

For in vivo  $\text{H}_2^{17}\text{O}$  imaging, a three-dimensional (3D) chemical shift imaging (CSI) sequence using weighted averaging of the k-space was applied (25,26). To obtain sufficient signal despite the short transverse relaxation time, the acquisition delay between the middle of the excitation pulse and the start of the acquisition was minimized to 484  $\mu$ s. Because a nonselective excitation pulse was used, the sensitive region in the field of view (FOV) of the 3D CSI sequence with free induction decay (FID) detection was determined by the profile of the surface coil. The raw data was Fourier transformed along the three spatial dimensions, resulting in localized FIDs, which were used to estimate  $T_2^*$  times. Subsequent Fourier transformation of the time domain resulted in localized spectra. The amplitudes of the phased water peaks were used to construct the final  $\text{H}_2^{17}\text{O}$  images.

### Phantom measurements

Reference relaxation time measurements on pure water with 20 mL of a 10%-enriched  $\text{H}_2^{17}\text{O}$  water phantom (Nukem GmbH, Germany) in a glass tube were performed at room temperature.

$T_1$  and  $T_2^*$  were measured as described below. A region of interest (ROI) in the center of the phantom was selected to avoid potential effects of superficial heating caused by the specific absorption rate (SAR) of RF-power

and to attenuate susceptibility effects in peripheral regions of the phantom.

### T<sub>1</sub> Relaxometry

Inversion recovery images were acquired with 20 inversion delays ranging from 0.5 ms to 45 ms in logarithmically increasing steps, using an adiabatic hyperbolic secant pulse (duration 0.5 ms, bandwidth 10 kHz) for inversion. Imaging parameters were: FOV 30 × 15.71 × 27 mm<sup>3</sup>; matrix size 21 × 11 × 9 (nominal voxel volume 6.12 μL); relaxation delay 45 ms (TR ≫ 5 T<sub>1</sub>); 50 μs pulse for excitation with a nominal flip angle of 90°. A total of 10,240 FIDs for each inversion step were acquired within 7–15 min (depending on the inversion time) with a maximum of 41 averages in the center of k-space. Spectroscopic data of 1250 points per FID was recorded within 10 ms with a bandwidth of 125 kHz (1321 ppm).

T<sub>1</sub> was determined by a mono-exponential fit of the intensity of selected voxels, using standard nonlinear least-squares algorithms (Curve Fitting Toolbox, MATLAB®, The Mathworks Inc., Natick, MA). The signal intensities for short inversion times were inverted while measurements within one standard deviation around the noise level were excluded before fitting.

The different tissue types were classified based on the proton images with reference to a rat brain atlas (27). At least two voxels between different tissue regions were not included along the outer boundary of ROIs to avoid partial volume contaminations due to the large voxel sizes in the H<sub>2</sub><sup>17</sup>O measurements. The T<sub>1</sub>-measurement covered a region of 0.9 to 1.5 cm in Z-direction (depending on the size of the rat), and thus most of each rat brain. Only voxels inside the brain with a sufficient spectral signal-to-noise ratio (SNR) > 12 were fitted, which resulted in a ratio of 56% brain versus 44% muscle voxels on average.

### T<sub>2</sub><sup>\*</sup> Relaxometry In Vivo

For T<sub>2</sub><sup>\*</sup> quantification, a short TR of 4.9 ms and a flip angle of 67.5°, close to the expected Ernst angle for the rat brain (FA = 64° for T<sub>1</sub> at 17.6T) (28) was used to sample 3D CSI with a FOV of 27.5 × 12.5 × 22 mm<sup>3</sup> and a 39 × 17 × 15 matrix (nominal voxel volume 0.76 μL). Three million FIDs (375 spectral points, bandwidth 100 kHz, 2811 averages in the center of k-space) were acquired within 4 h and 6 min from each rat.

The real part of the Fourier-transformed FID was frequency-corrected by adjusting the phase in the time-domain and fitted against a monoexponential decay. Voxel-wise estimated relaxation times were averaged over manually selected ROIs for each tissue type. Extrapolating the FIDs to zero excitation and acquisition delay resulted in T<sub>1</sub>-weighted images, after normalizing the intensity by correcting for B<sub>1</sub> inhomogeneity as determined in a phantom (29).

### T<sub>2</sub><sup>\*</sup> Relaxometry Postmortem

For high-resolution anatomical images and T<sub>2</sub><sup>\*</sup> quantification, a CSI-sequence without k-space weighting with a TR of 12 ms and a flip angle of 90° was used. To avoid RF-induced heating, the excitation pulse length was

Table 1

In Vivo T<sub>1</sub> Relaxation Times in ms of Natural Abundance H<sub>2</sub><sup>17</sup>O measured in the Rat Head

In vivo T <sub>1</sub>	Muscle tissue	Brain tissue
rat 1 (583 g)	5.25 ± 1.02 (324 μL)*	6.34 ± 0.67 (477 μL)*
rat 2 (415 g)	6.47 ± 1.15 (459 μL)	6.84 ± 0.62 (441 μL)
rat 3 (483 g)	5.86 ± 1.10 (275 μL)*	7.64 ± 0.52 (282 μL)*
rat 4 (580 g)	5.22 ± 0.84 (214 μL)*	6.58 ± 0.39 (398 μL)*
rat 5 (504 g) #1	5.45 ± 0.71 (312 μL)*	6.79 ± 0.40 (398 μL)*
rat 5 (504 g) #2	5.74 ± 0.70 (312 μL)*	6.83 ± 0.45 (398 μL)*
Pop. Mean [ms]	5.67 ± 1.12*	6.84 ± 0.67*

\*P < 0.05; significant population difference between muscle and brain tissue (paired t-test).

increased to 400 μs to compensate for the lack of the cooling effect of blood circulation.

A FOV of 27.5 × 12.5 × 25 mm<sup>3</sup> was sampled with a matrix of 41 × 19 × 25 voxels (voxel volume 0.44 μL). Approximately 2.5 million FIDs with 1000 points each and a spectral bandwidth of 100 kHz were acquired within 8 h and 18 min. Postprocessing for T<sub>2</sub><sup>\*</sup> was then performed as in vivo.

Relaxation times are given as mean value ± standard deviation over all voxels of the selected region, with the nominal volume of the ROI added in brackets, if applicable.

## RESULTS

### Phantom Measurements

Phantom measurements on the 10%-enriched H<sub>2</sub><sup>17</sup>O water sample resulted in a T<sub>1</sub> value of 6.24 ± 0.37 ms over the sensitive region of the surface coil, indicating stable T<sub>1</sub>-quantification and an efficient inversion over the entire FOV. In the same phantom, a T<sub>2</sub><sup>\*</sup> of 2.24 ± 0.09 ms was found.

### In Vivo Relaxation Times T<sub>1</sub>

T<sub>1</sub>-values for all rats are shown in Table 1. The T<sub>1</sub> of brain tissue (gray and white matter) was 6.84 ± 0.67 ms and significantly higher than that of muscle (5.67 ± 1.12 ms). Muscle tissue could be clearly differentiated from brain in image intensity (Fig. 1b) and relaxation times (Fig. 1c). No significant intracerebral contrast was observed at this resolution. A repeated measurement in the same rat in a different session demonstrated highly reproducible T<sub>1</sub>-values (Table 1; rat 5).

### In Vivo and Postmortem Relaxation Times T<sub>2</sub><sup>\*</sup>

In the high-resolution H<sub>2</sub><sup>17</sup>O images, intracortical structures were consistently visible due to the decreased image intensity of white matter (i.e., corpus callosum and optic radiation) against cortical gray matter in all animals, both in vivo and postmortem. T<sub>2</sub><sup>\*</sup> of brain tissue had a mean value of 1.77 ± 0.04 ms and did not vary significantly between white and gray matter (Fig. 2c; Table 2a), while being significantly increased in cerebrospinal fluid (CSF), e.g., in the ventricles (in vivo: 2.21 ± 0.20 ms; postmortem: 2.30 ± 0.16 ms, Table 2). A T<sub>2</sub><sup>\*</sup> of ~4 ms was measured in the eyes, too small in volume for quantitative evaluation. No significant deviation from a single exponential relaxation behavior was



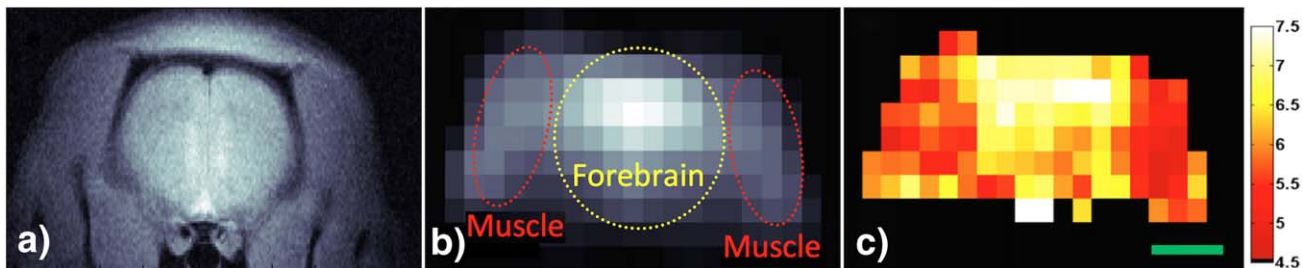


FIG. 1. Longitudinal relaxation time  $T_1$ . **a**: Proton FLASH image of a coronal slice in the rat head in vivo. **b**: Signal intensity of the  $H_2^{17}O$  spectral peak at natural abundance concentration of the same slice with delineations of forebrain (yellow) and surrounding muscle (red). **c**: Corresponding slice of  $H_2^{17}O$  in vivo  $T_1$  relaxation times with clear differences between brain and muscle tissue. A 5-mm scale bar (green) is shown for the colocalized slices.

apparent in the signal time course of the localized FIDs as used for the  $T_2^*$  estimations (Fig. 3a). Extrapolation to  $TE = 0$  yielded combined spin density- and  $T_1$ -weighted images (Fig. 3c), showing highest signal intensity in muscle tissue and lowest in CSF and ventricles, confirming the lower  $T_1$  in muscle than in brain (Fig. 1c).

Rats undergoing  $^{17}O_2$ -inhalations which were used for the postmortem  $T_2^*$  measurements showed a global increase of the image intensity by a factor of 4 to 5 compared with the natural abundance  $H_2^{17}O$  signal, due to significant metabolic conversion of  $^{17}O_2$  to  $H_2^{17}O$ . No intracortical difference from metabolic rates (e.g., between white and gray matter) was detectable, due to the long duration of anatomical acquisition (i.e., diffusion of locally elevated  $H_2^{17}O$  concentration in equilibrium due to exhaustion of concentration gradients). Postmortem  $T_2^*$  values in muscle, GM and WM tissue were slightly increased compared with in vivo data, whereas the difference between brain tissue and CSF remained similar (Table 2b). This was accompanied by a generally decreased image contrast in the postmortem measurements, especially in regions with high water content (Fig. 4).

## DISCUSSION

In this study, we have taken advantage of the increased signal at ultra-high magnetic field to quantify  $H_2^{17}O$

relaxation times in the rat brain in vivo, for the first time differentiating different tissue types with relatively high spatial resolution. The results of this study should be beneficial to design future experiments, but also to assess the results of previous quantitative examinations. While a spectroscopic technique was used in all acquisitions to reach high sensitivity for the rats with and without  $^{17}O_2$  inhalation, no metabolite peaks except the water signal were observable in our spectra (30–34) because of the limited excitation bandwidth in relation to the large range of chemical shifts with molecules containing  $^{17}O$ .

## Phantom Studies

Our phantom measurements are in relative agreement with results from previous nonlocalized measurements which reported values of  $T_1$  around 6–8 ms for enriched pure water (35) at physiological temperatures. Temperature controlled  $T_2^*$  measurements using continuous wave NMR reported  $T_2^*$  up to 4 ms at enriched  $H_2^{17}O$  concentrations (36,37) in contrast to our significantly lower values, which could be a result of the higher achievable  $B_0$  field-homogeneity in nonlocalized NMR spectroscopy, of better temperature control in the nonlocalized measurements in NMR spectrometers and higher purity from paramagnetic ions of the glassware being

Table 2  
In Vivo and Postmortem  $T_2^*$  Relaxation Times in ms of  $H_2^{17}O$  measured in the Rat Head

a) In vivo $T_2^*$ (natural abundance)				
	Muscle tissue	WM	GM	CSF
rat A (567 g)	1.03 ± 0.18 (309 $\mu$ L)	1.86 ± 0.15 (24 $\mu$ L)	1.83 ± 0.13 (111 $\mu$ L)	2.56 ± 0.23 (38 $\mu$ L)
rat B (499 g)	1.07 ± 0.18 (296 $\mu$ L)	1.78 ± 0.09 (29 $\mu$ L)	1.75 ± 0.08 (126 $\mu$ L)	2.19 ± 0.12 (27 $\mu$ L)
rat C (503 g)	1.08 ± 0.17 (301 $\mu$ L)	1.75 ± 0.11 (25 $\mu$ L)	1.73 ± 0.10 (128 $\mu$ L)	2.13 ± 0.16 (28 $\mu$ L)
rat D (526 g)	1.15 ± 0.16 (361 $\mu$ L)	1.71 ± 0.10 (20 $\mu$ L)	1.77 ± 0.08 (148 $\mu$ L)	2.14 ± 0.11 (27 $\mu$ L)
rat E (552 g)	1.27 ± 0.15 (315 $\mu$ L)	1.77 ± 0.08 (14 $\mu$ L)	1.76 ± 0.15 (100 $\mu$ L)	2.05 ± 0.11 (10 $\mu$ L)
Pop. Mean [ms]	1.12 ± 0.09*	1.77 ± 0.06	1.77 ± 0.04	2.21 ± 0.20*
b) Postmortem $T_2^*$ (enriched ~4–5 times of natural abundance)				
rat F (250 g)	1.30 ± 0.21 (232 $\mu$ L)	1.81 ± 0.17 (33 $\mu$ L)	1.83 ± 0.15 (199 $\mu$ L)	2.27 ± 0.35 (14 $\mu$ L)
rat G (430 g)	1.28 ± 0.16 (239 $\mu$ L)	1.84 ± 0.21 (25 $\mu$ L)	1.85 ± 0.17 (182 $\mu$ L)	2.26 ± 0.39 (9 $\mu$ L)
rat H (540 g)	1.29 ± 0.25 (276 $\mu$ L)	1.91 ± 0.33 (24 $\mu$ L)	1.92 ± 0.21 (102 $\mu$ L)	2.14 ± 0.44 (8 $\mu$ L)
rat I (300 g)	1.20 ± 0.16 (207 $\mu$ L)	1.87 ± 0.16 (22 $\mu$ L)	1.89 ± 0.12 (175 $\mu$ L)	2.53 ± 0.35 (25 $\mu$ L)
Pop. Mean [ms]	1.27 ± 0.05*	1.86 ± 0.04	1.87 ± 0.04	2.30 ± 0.16*

\* $P < 0.01$  significant population difference between muscle/CSF and brain tissue (paired t-test).

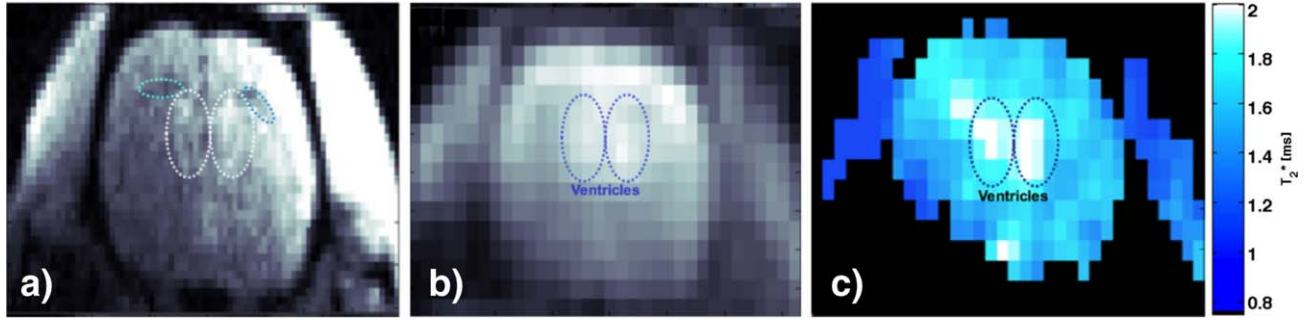


FIG. 2. Apparent in vivo transverse relaxation time  $T_2^*$ . **a:** Coronal FLASH proton image. **b:** CSI anatomical contrast of rat tissues in image intensity of  $\text{H}_2^{17}\text{O}$  with intra-cortical differentiations of white matter and ventricles. **c:**  $T_2^*$  estimates of  $\text{H}_2^{17}\text{O}$  with pronounced differentiation between brain parenchyma (gray matter, white matter) and cerebrospinal fluid in the ventricles.

used (35). Furthermore, a fair amount of one-dimensional experiments in the 1970s accumulated in vitro estimations of muscle tissue  $T_1$  and  $T_2$  (37–39). In contrast to proton MRI,  $T_2^*$  contrast changes close to air-sample borders were not observed in any images nor with varying voxel sizes as the frequency of  $\text{H}_2^{17}\text{O}$  at

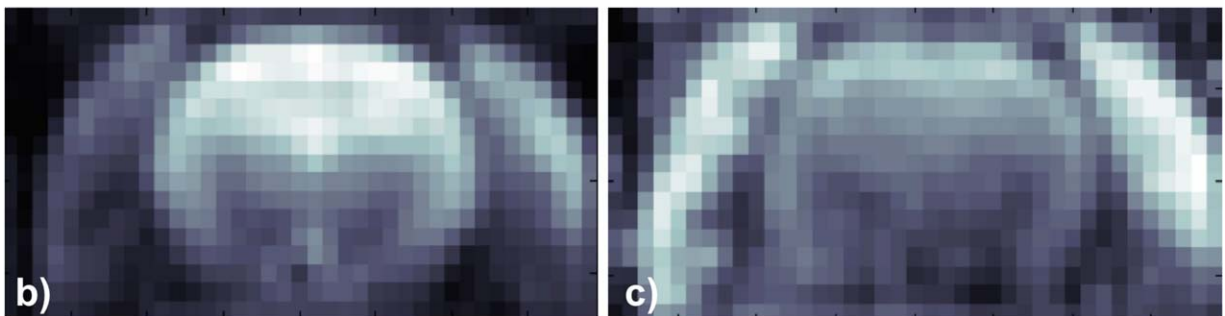
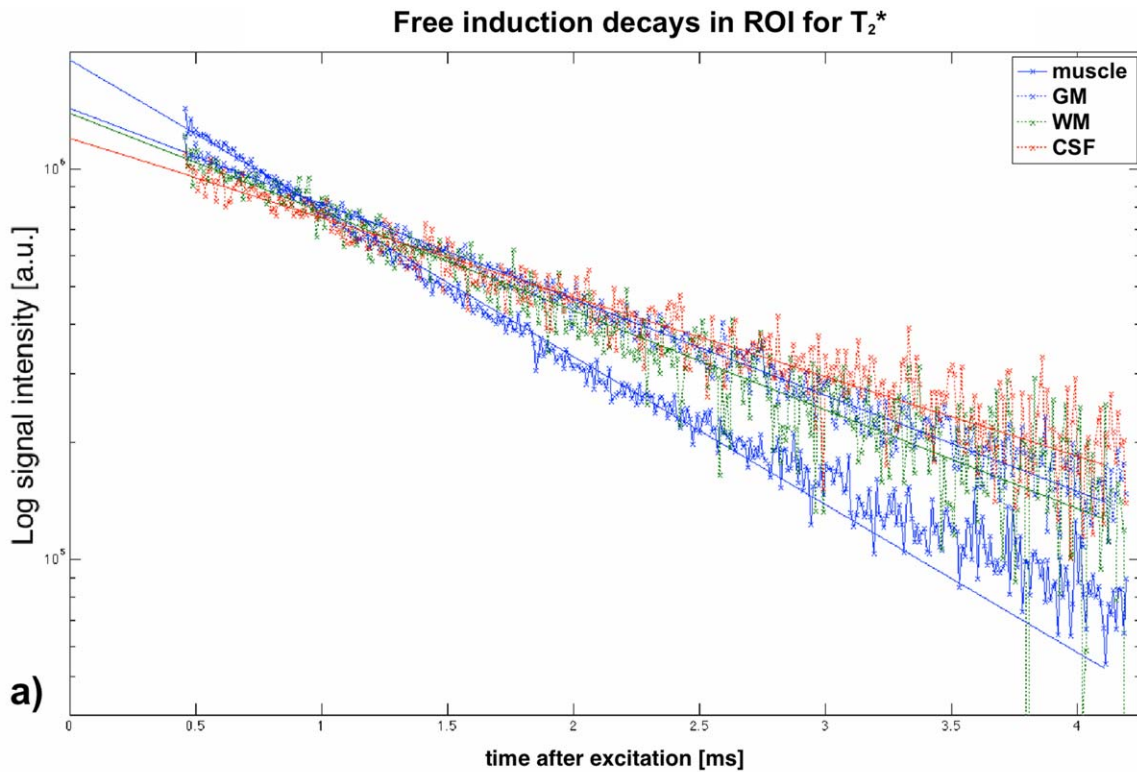


FIG. 3. **a:** Semi-logarithmic plot of the MR-signal (real part) averaged over ROIs of one rat from the in vivo measurements with fitted  $T_2^*$  relaxation curves for four tissue types. Free induction decays of different tissue types show high linearity without pronounced modulations. **b:** Reconstructed image of muscle and brain tissue based on the spectral peak of  $\text{H}_2^{17}\text{O}$ . **c:**  $T_1$ -weighted image generated by extrapolation of the FIDs (see text) to an acquisition delay of  $t = 0$  for the same coronal slice as shown in b).

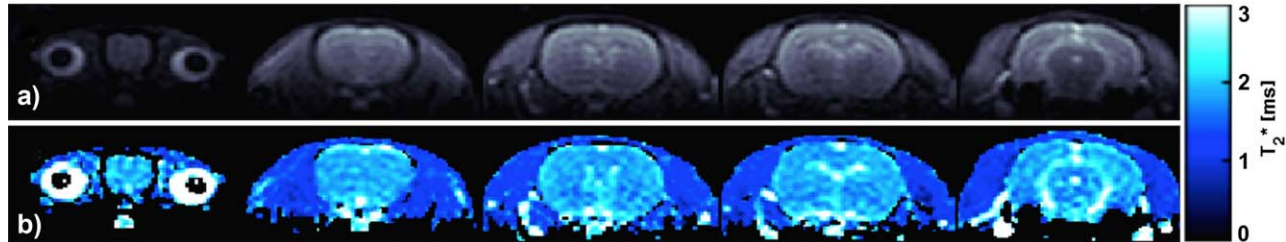


FIG. 4. Apparent postmortem transverse relaxation time  $T_2^*$ . **a:** Postmortem images with increased  $H_2^{17}O$  concentration due to metabolic enrichment after inhalation of  $^{17}O_2$  gas. Olfactory bulb, forebrain, thalamus and cortex are shown with clear contrast between white matter (corpus callosum) and ventricles. **b:**  $T_2^*$  map of the same slices as in a) exhibiting slightly increased relaxation times in cortical fluids of the ventricles.

16.4 Tesla used in this study translates to a larmor frequency of 2.2T for protons. In addition, peak intensity changed by less than 15% comparing unlocalized spectra with and without  $B_0$ -shimming, (i.e., with a smaller change in  $T_2^*$ ).

#### Tissue Specific Relaxation and Anatomical Contrast

The differences between muscle and brain tissue  $H_2^{17}O$  relaxation times have similar proportions as those known from proton measurements (40–42) and cannot be explained from temperature difference alone. However, intracortical differences are much less pronounced than for proton MRI due to the limited influence of the intermolecular environment on the field-independent relaxation of  $H_2^{17}O$  (43). The lack of differentiable  $T_2^*$  relaxation times between gray and white matter and the clear separation by low intensity of white matter structures indicate a difference in either  $T_1$  or in  $H_2^{17}O$  concentration (Fig. 3c). The latter could be due to the variations in water content of more than 15% between gray and white matter as found in the human brain (44) and could affect natural abundance  $H_2^{17}O$  calibration (45).

The higher intersubject variability of the  $T_1$  measurements compared with the  $T_2^*$  data (Table 1 versus 2) is mainly due the lower sensitivity of the inversion recovery sequence compared with the simple FID acquisition used to determine  $T_2^*$ . In part, choosing larger voxels for the  $T_1$  measurements compensated for this. However, the low sensitivity of  $H_2^{17}O$  in general requires multiple repetitions for averaging the signal both in  $T_1$  and  $T_2^*$  measurements, which is possible due to the generally short  $T_1$  in contrast to more slowly relaxing nuclei like protons.

Most previous investigations on oxygen consumption were limited in localization because of the low sensitivity of the  $^{17}O$  nucleus, thus neglecting potential contributions from nonbrain tissue. In animal studies using surface coils for detection, the signal from muscles could cause significant confound, owing to the property of small coils to overemphasize superficial (i.e., proximal) anatomical structures. However, our results indicate that the low  $T_2^*$  values in muscle tissues presented here assure that the signal contribution from this tissue type will play a less dominating role when sufficiently long TE is used.

While the still relatively large voxels sizes may lead to residual partial volume effects, further increasing the

spatial resolution, which is constricted by the low SNR of the measurements, will be challenging. In our experiments, we have chosen a CSI sequence because of its high SNR efficiency and its short acquisition delay. The necessary SNR gain for improving the resolution might be possible by further reducing the signal loss due to transverse relaxation by using imaging sequences with ultrashort or even zero TE.

#### Comparison of In Vivo and Postmortem

Differences between in vivo and postmortem relaxation times reflect well known effects from previous studies (e.g., due to cell rupture, increased diffusion between intra- and extracellular space and differences in temperature) (23,46). Temperature changes due to the rapid succession of RF-pulses can distort the resulting relaxation times especially in the postmortem measurements, where the body temperature is lower and no perfusion dissipates the heat. While SAR was reduced in these measurements by increasing the pulse duration, regions close to the coil, especially in the skeletal muscles, may heat up, explaining the slightly higher  $T_2^*$  in those regions. In the ex vivo measurements, it was possible to reach

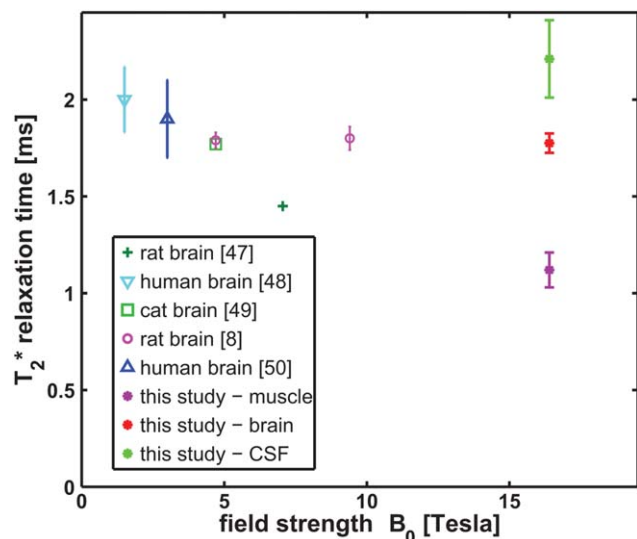


FIG. 5. Literature comparison of  $T_2^*$  relaxation times at different  $B_0$  field strength of in vivo brain tissue. Span of error bar equals standard deviation, if available.



even higher spatial resolutions due to the H<sub>2</sub><sup>17</sup>O enrichment and the long measurement times.

### <sup>17</sup>O Relaxation Times across Field Strength

In contrast to proton MRI, the longitudinal and transverse relaxation times of H<sub>2</sub><sup>17</sup>O are expected to be independent of field strength. Figure 5 shows an overview of previously published T<sub>2</sub>\*-values at different field strengths (8,47–50), including the results from the current study. All brain tissue T<sub>2</sub>\* values are remarkably close around 1.8 ms, because the tissue is maintained in an auto-regulatory regime of temperature and pH, some of the most dominant factors for transverse H<sub>2</sub><sup>17</sup>O relaxation. This data is consistent with the assumption of field-independent relaxation times, albeit a slight tendency to shorter T<sub>2</sub>\* (Fig. 5) might be visible in muscle tissue.

Differences to previous studies might be due to differences in temperature, which may result from convection effects or the different temperature controlling procedures. They can also be caused by RF-heating in in vivo experiments, because the short T<sub>1</sub> and T<sub>2</sub>\* and the necessity of rapid averaging increase SAR. For <sup>17</sup>O experiments at high fields, SAR monitoring and reduction is critical especially with humans. Thus, in the postmortem study we used longer RF excitation pulses to avoid excessive heating at the expense of signal loss due to relaxation during the pulse. Further discrepancies between the current study and previous experiments may be due to differences in the sampled tissue, because most previous in vivo data were acquired with limited spatial localization and might represent a different mixture of tissue types.

### CONCLUSIONS

The use of ultrahigh field strength made it possible to determine tissue specific <sup>17</sup>O relaxation parameters. These might be used for sequence optimization (13,51) and analysis of data acquired using short TR during inhalation of enriched <sup>17</sup>O<sub>2</sub> gas for the determination of CMRO<sub>2</sub>. Due to the apparent field independence of the <sup>17</sup>O relaxation times, the results can easily be transferred to lower field strengths, as long as tissue and temperature differences are taken into account.

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