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Carbon bed post-plasma to enhance the CO₂ conversion and remove O₂ from the product stream

Fanny Girard-Sahun^{a,*}, Omar Biondo^{a,b,*}, Georgi Trenchev^a, Gerard van Rooij^{b,c}, Annemie Bogaerts^{a,*}

^a Research Group PLASMANT, Department of Chemistry, University of Antwerp, Universiteitsplein 1, Wilrijk B-2610, Belgium

^b DIFFER, 5612AJ Eindhoven, The Netherlands

^c Faculty of Science and Engineering, Maastricht University, 6229 GS Maastricht, The Netherlands

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ABSTRACT

CO₂ conversion by plasma technology is gaining increasing interest. We present a carbon (charcoal) bed placed after a Gliding Arc Plasmatron (GAP) reactor, to enhance the CO₂ conversion, promote O/O₂ removal and increase the CO fraction in the exhaust mixture. By means of an innovative (silo) system, the carbon is constantly supplied, to avoid carbon depletion upon reaction with O/O₂. Using this carbon bed, the CO₂ conversion is enhanced by almost a factor of two (from 7.6 to 12.6%), while the CO concentration even increases by a factor of three (from 7.2 to 21.9%), and O₂ is completely removed from the exhaust mixture. Moreover, the energy efficiency of the conversion process drastically increases from 27.9 to 45.4%, and the energy cost significantly drops from 41.9 to 25.4 kJ.L⁻¹. We also present the temperature as a function of distance from the reactor outlet, as well as the CO₂, CO and O₂ concentrations and the temperature in the carbon bed as a function of time, which is important for understanding the underlying mechanisms. Indeed, these time-resolved measurements reveal that the initial enhancements in CO₂ conversion and in CO concentration are not maintained in our current setup. Therefore, we present a model to study the gasification of carbon with different feed gases (i.e., O₂, CO and CO₂ separately), from which we can conclude that the oxygen coverage at the surface plays a key role in determining the product composition and the rate of carbon consumption. Indeed, our model insights indicate that the drop in CO₂ conversion and in CO concentration after a few minutes is attributed to deactivation of the carbon bed, due to rapid formation of oxygen complexes at the surface.

1. Introduction

Due to the high CO₂ levels in the atmosphere (more than 410 ppm [1]), we need to find urgent solutions for reducing CO₂ emissions in our fight against global warming. Since 2015, with the Paris agreement, there have been multiple efforts towards the implementation of low-carbon energy production, in order to pursue the net-zero target, in at least 20 countries in the world. A net-zero emissions energy system does not add any CO₂ to the atmosphere [2]. There are different ways of approaching the issue of reducing CO₂ emissions. For instance, H₂ has the highest possible energy density and its use features CO₂-free emissions. However, hydrogen-based energy production is limited by the storage capacity of the available facilities, due to the very low

volumetric density of hydrogen. Such a drawback can be overcome with the development of the offshore geological storage of hydrogen (OGSH) in sub-sea reservoirs [3]. Alternatively, CO₂ can be either captured and stored (known as CCS) or captured and utilized into valuable chemicals, useful for a wide range of industrial applications. This approach, known as “CCU - CO₂ Capture and Utilization” is very promising, as CO₂ becomes then a renewable resource of energy for various applications. Especially, CO₂ utilization as feedstock for fuels represents a major route for decarbonization [1].

Direct CO₂ splitting (equation (1)) is highly endothermic and requires a lot of energy in thermo-catalytic processes, i.e. temperatures of at least 2000 K at atmospheric pressure [4]. In this context, plasma technology is a very promising candidate to face this challenge [4]. Due to its unique and highly-reactive composition (high-energy free

* Corresponding authors at: Research Group PLASMANT, Department of Chemistry, University of Antwerp, Universiteitsplein 1, Wilrijk B-2610, Belgium.

E-mail addresses: fanny.girard-sahun@uantwerpen.be (F. Girard-Sahun), omar.biondo@uantwerpen.be (O. Biondo), annemie.bogaerts@uantwerpen.be (A. Bogaerts).

¹ Shared first author.

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Nomenclature*Latin letters*

A	cross-sectional area (cm ²)
a	stoichiometric coefficient
B	active site density (mol.m ⁻²)
b	fitting parameter
C _{surf}	active site density (cm ⁻³)
C(s)	active carbon site
C(s)-O	oxygen atom chemisorbed at an active carbon site
EC	energy cost (kJ.L ⁻¹)
h	Planck constant (J.s)
ΔH ^o _R	standard reaction enthalpy (kJ.mol ⁻¹)
k	reaction rate coefficient
k _B	Boltzmann constant (J.K ⁻¹)
m	mass of one molecule (kg)
N _A	Avogadro number (mol ⁻¹)
n	number density
P	power (kW)
P	pressure (Pa)
Q _{in}	input flow rate (mol.min ⁻¹)
Q _{out}	output flow rate (mol.min ⁻¹)
R	ideal gas constant (L.atm.K ⁻¹ .mol ⁻¹)
R	reaction rate
SEI	specific energy input (kJ.mol ⁻¹)
T	temperature (K)
Td	Townsend; 1 Td = 10 ²¹ V.m ⁻²
V _m	molar volume at atmospheric pressure and room temperature (L.mol ⁻¹)
V	volume (cm ³)
X	conversion (%)
x	position in the simulated reactor

Greek letters

π	pi number
∏	product

Abbreviations

APS	Advanced Plasma Solutions
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BET	Brunauer-Emmett-Teller
CAPEX	Capital expenditure
CCS	CO ₂ capture and storage
CCU	CO ₂ capture and utilization
DBD	dielectric barrier discharge
DFT	density functional theory
eV	electronvolt
GAP	gliding arc plasmatron
GC	gas chromatography
MW	microwave
NDIR	non-dispersive infrared spectroscopy
OGSH	offshore geological storage of hydrogen
OPEX	operating expenses
RBR	reverse Boudouard reaction
RGA	residual gas analysis
SEM	scanning electron microscopy
TCD	thermal conductivity detector
TGA	thermogravimetric analysis
TGA-MS	thermogravimetric analysis coupled with mass spectrometry
TSA	total surface area
OD	zero-dimensional

Superscripts

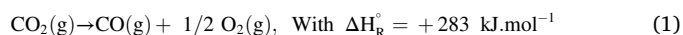
L	left-hand side
R	right-hand side

Subscripts

bed	carbon bed
corr	corrected for the gas expansion
0	Initial
j	reaction j
in	Input
l	reactants of reaction j
out	Output
s	Species
surf	Surface

electrons, excited molecules/atoms, radicals and ions, as well as photons), plasma offers a route to convert CO₂ with high energy efficiency and operation flexibility, and at mild conditions (ambient pressure and temperature), as demonstrated in numerous studies [5]. For instance, pure CO₂ dissociation was investigated with dielectric barrier discharges (DBD), which are one of most common types of non-thermal plasmas for CO₂ conversion [4,6,7]. Notwithstanding this, DBDs provide low energy efficiency (around 10%) [4]. Under typical conditions for DBDs, i.e. reduced electric field of 200 Td (Townsend; 1 Td = 10⁻²¹ V m²) [4] and electron temperature higher than 5 eV [8], only ca. 10% of the electron energy is deposited into the CO₂ vibrational degrees of freedom, which are known to promote dissociation through the most efficient pathway [4]. Therefore, intensive research has been conducted on the applicability of so-called “warm” plasmas, which are plasmas with a high degree of equilibrium between vibrational and translational degrees of freedom. The particular conditions of operation of these plasmas allow combining advantages of both thermal (high gas temperature) and non-thermal (vibrational excitation) systems [4,9]. With this aim, “warm” plasmas such as microwave (MW) plasmas [10-12], (2D) gliding arc plasma [13,14] or (3D) gliding arc plasmatron at atmospheric pressure [14-16], were investigated for more efficient CO₂ conversion, providing promising results. Recently, an atmospheric pressure glow discharge [15] was also proposed as an interesting technology for CO₂

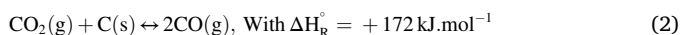
dissociation, offering a conversion degree of up to 12% and energy efficiencies reaching 30%. Additionally, the conversion of CO₂ mixed with N₂, which is a typical impurity of waste streams and which can promote CO₂ vibrational excitation through efficient vibrational exchanges [16,17], was also studied in gliding arc reactors [18,19] and in a DBD reactor [20]. The dissociation of CO₂ leads to the formation of CO and an oxygen atom, and is a highly endothermic reaction, requiring a considerable amount of energy.



The CO produced can be used in e.g., the Fischer-Tropsch process to produce liquid hydrocarbons. Currently, plasma technology for CO₂ conversion is still at low technological readiness level (TRL 3-4; lab scale). However, plasma is a well-established technology for lighting, surface cleaning, etching, film deposition and polymerization at industrial level [21,22]. Moreover, non-thermal plasma stands out as a promising alternative to the widely used oxidation processes for water and flue gas treatment [23] and can be a suitable tool for nanomaterial processing, analytical chemistry, sterilization, disinfection, medicine and food safety [24].

More recently, another possible route for CO₂ conversion using biochar as a solid reactant is being investigated, also called CO₂ gasification of biochar [25]. Biochar is obtained from pyrolysis of biomass and

it is a quite cheap renewable energy source [26]. With this in mind, and provided that the system is powered by electricity coming from renewable sources, we consider CO₂ gasification of biochars a promising candidate to bolster the efforts towards a net-zero energy system. Furthermore, gasification of coal is a well-known process for the production of CO and syngas. We believe that the combination of these technologies can bring a novel, highly efficient solution to the market. The reaction of CO₂ with carbon is known as the Reverse Boudouard Reaction (RBR) (equation (2)) [27]:



At atmospheric pressure, the RBR also requires high temperatures (>700 °C) [27]. Several gasification heat sources to favour the RBR rate have been studied, e.g. conventional convective heating [27], solar-driven gasification [27] and microwave-driven CO₂ gasification [28–30]. Especially MW heating [28–30] allowed to decrease the activation energy for CO₂ gasification by 60–70% and to reach energy efficiencies of 45% at laboratory scale [31]. Recently, a few groups investigated the RBR with solid carbon driven by different kinds of plasmas and showed very interesting results [32–34].

Overall, one should consider that O₂-free exhaust gas is needed from CO₂ splitting (i.e., ideally pure CO) to be directly used in industrial applications, like the Fischer-Tropsch process as mentioned above. Indeed, CO₂ splitting results in CO and O, and the latter will (1) recombine to form O₂, but also (2) be responsible for unwanted back-reactions (CO + O) to produce again CO₂ [35], lowering the net CO₂ conversion and energy efficiency. In addition, also O₂ can give rise to back-reactions with CO into CO₂. Therefore, both O and O₂ are undesirable species that need to be either physically or chemically quenched to avoid reverse reactions, reduce separation costs downstream and increase the energy efficiency [6,36,37]. Only a few methods are described in literature, such as the use of hollow fiber membranes permeable to O/O₂, which is then removed using a sweep gas [38]. A promising route, also mentioned above, is to use solid carbon after the plasma reactor, which can be oxidized by O and/or O₂ to produce CO, allowing (i) to consume the undesirable O/O₂ and (ii) to enrich the exhaust in CO, beneficial for the applications.

In this work, we investigate the conversion of CO₂ by a Gliding Arc Plasmatron (GAP) in combination with a carbon bed after the plasma reactor, to promote O/O₂ removal and increase the CO fraction in the exhaust mixture. Although the plasma method requires quite some energy, it uses electricity, and is quickly switched on/off, so it is very suitable for combination with renewable electricity. Moreover, we demonstrated already high energy efficiency for plasma-based CO₂ conversion in this GAP [39]. The combination with carbon bed will further improve the performance. We present an innovative design, in which the carbon is constantly supplied over time by a silo system, to avoid carbon depletion upon reaction with O/O₂. The carbon bed is put in direct contact with the afterglow of the gliding arc plasma produced in pure CO₂ at atmospheric pressure, for maximum effect. During the process, the temperature is recorded as a function of time directly in the carbon bed, providing useful information for a better understanding of the underlying mechanisms. Furthermore, the exhaust gas composition is analysed by a combination of gas chromatography and optical sensors. The performance of the GAP with carbon bed is determined and compared with the benchmark (i.e. without carbon bed).

Although it has been experimentally proven that carbonaceous materials are effective options to increase the selectivity towards CO in the CO₂ splitting process, more insight is needed to demonstrate its full potential and identify the relevant underlying mechanisms. In this respect, modelling is a very useful tool to unravel chemical pathways especially when these are inaccessible experimentally. Zero-dimensional (0D) chemical kinetic models, also referred to as global models, are the method of choice to study the contributions of different chemical reactions occurring in plasma ignited in complex gas mixtures [40].

Therefore, in addition to our experimental work, we examine here the gasification process occurring when a reactive gas mixture produced by a plasma comes in contact with a carbon bed. This analysis involves the construction and the validation of a 0D chemical kinetics model against experiments. The validation of a model is a crucial aspect and requires particular experimental conditions, in order to minimize the number of assumptions under which the predictions are valid. With this in mind, we have chosen to validate our model with the gasification experiments presented by Panerai *et al.* [41], under pure O₂ and CO₂ atmospheres, respectively, in order to get insights into the reaction pathways for carbon oxidation in the presence of O₂, CO₂ and CO and the deactivation of the carbon bed, with consequent decrease in selectivity towards the latter. This model greatly helps us to explain our experimental results, as will be described in detail in this paper. To resume, the innovations of this work, compared to [32–34], are: i) the use of a silo system, to avoid carbon depletion due to gasification; ii) time-resolved temperature measurements in the carbon bed; iii) a combination of gas chromatography and optical sensors to determine the exhaust composition as a function of time; iv) and especially coupling our experiments with detailed kinetic modelling, to provide detailed insights into the underlying chemistry and the deactivation of the carbon bed.

2. Experimental section

2.1. Experimental setup

We performed our experiments in a gliding arc plasmatron (GAP) originally developed by Nunnally *et al.* [42] and previously described by our group [39]. We used pure CO₂ as reactant gas in the plasma reactor (Air Liquide, purity 99.995%; flow rate kept constant at 10 L_n·min⁻¹). The flow rate was adjusted with a Bronkhorst mass flow controller. The cathode and anode (stainless steel) were connected to an AC current source-type power supply (AFS) and to the ground, respectively. Electrode degradation (e.g., due to ablation) was very limited at the conditions under study. Indeed, this AC power supply reduces the thermal stress on the developing cathode spots. Furthermore, we utilize a strongly turbulent internal gas flow, which efficiently cools down the electrodes. Voltage and current were measured by using a high-voltage probe (Cal Test Electronics CT4028) and a current sense resistor (2 Ω) connected to a two-channel digital oscilloscope (Keysight DSOX1102G). Based on the product of measured voltage and current, we calculated the power injected in the plasma. The plasma arc was formed between both electrodes, as previously described [39], and can also partly leave the reactor through the anode (reactor outlet) [43], forming an afterglow in contact with the carbon bed, placed at the anode exhaust.

The general scheme, including the GAP plasma reactor, the carbon bed and the silo, is illustrated in Figure 1A, and a picture is given in Fig. 1B.

The carbon basket is directly attached to the anode. The silo containing spare carbon is attached to the basket through the exhaust reactor body. A metallic mesh is used at the anode exit to keep the carbon pellets out of the plasma reactor. Two dedicated ports allow to measure the temperature inside the carbon bed, at 15 and 35 mm from the anode exit; see Fig. 2 below.

The reactor outlet (length 26 mm and inner diameter 7.07 mm, see pictures in Supporting Information (SI): Figure S1) was designed so that the contact of the plasma afterglow and the carbon bed is maximized: the aperture on the basket side is wider (see Figure S1(C) in SI), so that the carbon can get closer to the arc, which improves the contact between active plasma species and the carbon, and also takes benefit of the high temperature of the arc (3000 K according to computer simulations) [44].

The basket was filled with ca. 5 g of carbon pellets for each plasma run. We used two different types of commercial carbon materials of different quality and particle size: activated charcoal Norit PK 1–3 (Sigma Aldrich, 1–3 mm typical pellet size, from peat, steam activated; named “charcoal 1” in the figures below) and untreated activated

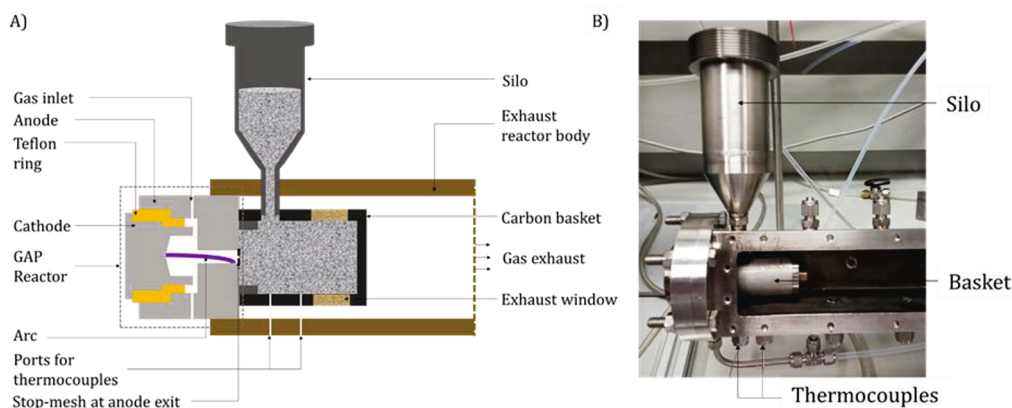


Fig. 1. Scheme (A) and picture (B) of the GAP plasma reactor setup, including the carbon basket and silo. Note: The scheme is not at scale.

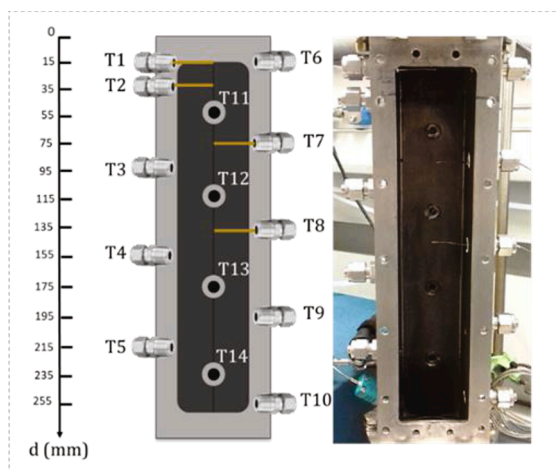


Fig. 2. Scheme (left) and picture (right) of the post-plasma reactor tube with several ports for the K-type thermocouples (named Tx, with x = 1 to 14). The “zero” position corresponds to the reactor outlet (anode exit position). T1 and T2 (at 15 and 35 mm) measure the temperature inside the carbon bed; cf. Fig. 1).

charcoal C3014 (Sigma Aldrich, typical pellet size of 250–850 μm ; named “charcoal 2”). Note that there is a very wide variety of commercial carbon samples available in the market. The particle size is usually the main (even sometimes the only one) parameter given by carbon suppliers. Therefore, it is the easiest way to investigate differences among the wide range of carbon samples available. In future work, we plan to also compare carbon samples with the same particle size but different specific surface area and morphology, if available. The surface area of charcoal 1 and charcoal 2, as received, was estimated by N_2 -Brunauer-Emmett-Teller (BET) measurements to be 705 ± 71 and $703 \pm 70 \text{ m}^2 \cdot \text{g}^{-1}$, respectively (see section 2.2. for details on the technique). The surface area of charcoal 1 was also measured after 7 min of treatment, in order to estimate the surface variation resulting from gasification reactions. The morphology of the charcoals as received, as well as after the gasification reactions, was investigated by Scanning Electron Microscopy (SEM-see section 2.2.). The collected images are reported in the Supporting Information (SI), in Figure S2.

A silo (stainless steel) containing fresh carbon was inserted inside the basket through the post-plasma reactor tube, to provide a continuous source of carbon during the reaction and avoid carbon depletion upon reaction with O/O_2 (within the limits of the carbon volume contained in the silo, i.e., approx. 200 cm^3 , see Fig. 1A). It can be refilled with carbon through an opening on its top (Fig. 1B). A mesh (stainless steel) was added at the anode exit to hold the carbon pellets inside the basket and

avoid disturbing the arc inside the plasma reactor by creating a short circuit. In addition, the temperature in the carbon bed during operation was measured by means of two thermocouple ports along the basket (see Fig. 1A-B).

Filters were placed on the exhaust line to make sure that the charcoal/ash particles did not enter the gas analytic equipment. The exhaust gas composition (CO_2 , CO and O_2) was analysed online by NDIR (Non-Dispersive Infrared Spectroscopy, for CO_2 and CO) and by an optical oxygen sensor (for O_2). Besides, the exhaust gas mixture was also analysed by gas chromatography (GC); see next section.

The post-plasma reactor tube was designed specifically for the purpose of these experiments, i.e., (1) to accommodate the carbon basket combined with the silo, and (2) to measure the temperature in the post-plasma region as a function of distance from the reactor outlet, thanks to several ports for thermocouples. Indeed, to determine the best position of the carbon bed after the plasma at our operating conditions, we first measured the temperature profile in the post-plasma reactor tube with K-type thermocouples with four-channels type data loggers (Thermosense) in a free-reactor (i.e., without carbon bed; see Fig. 2).

2.2. Carbon materials characterization techniques

The carbon materials (namely charcoal 1 and 2) used in this study were characterized by Brunauer-Emmett-Teller (BET) analysis, Scanning Electron Microscopy (SEM) and Thermogravimetric Analysis coupled with Mass Spectrometry (TGA-MS).

The specific surface area of charcoal 1 and 2 was calculated by the BET method. Prior to the measurements, the samples were degassed for 16 h under vacuum at $200 \text{ }^\circ\text{C}$. The sorption was then performed under N_2 atmosphere using a Quantachrome Quadrasorb SI.

For SEM analysis, an FEI Quanta 250 FEG scanning electron microscope was used. It was operated at 20 kV and the secondary electron detector was used to generate the images. The working distance (distance between the sample and the polepiece, i.e. the last lens of the imaging system; “WD”) is given in the bar below every image and was always between 8 and 10 mm. To prepare the sample, a chunk of the charcoal was attached to the aluminium SEM stub using double sided carbon tape. No further treatment of the sample was needed (like coating or crushing).

Thermogravimetric analysis (TGA) was used for studying the thermal properties of the activated charcoal, using a Mettler Toledo TGA-DSC 3+. The measurement was performed under a continuous flow of argon or oxygen i.e. 100 mL/min (80 mL/min protective gas to protect the balance + 20 mL/min purge gas). The sample was first flushed 45 min at $30 \text{ }^\circ\text{C}$ to remove all the air from the oven, and then heated up from 30 to $1000 \text{ }^\circ\text{C}$ with a heating rate of $10 \text{ }^\circ\text{C}/\text{min}$. An alumina crucible of $150 \mu\text{L}$ was used. TGA was associated with a mass spectrometer to analyse the products from charcoal pyrolysis. The detection of the gaseous products

was performed with a HPR-20 R&D (Hidden) mass spectrometer, controlled by EGAsoft software. The following settings were used: an electron voltage of 70 eV and an emission current of 400 mA. An electron multiplier (SEM) detector was used to scan m/z from 1 to 19, from 21 to 39 and from 42 to 120.

2.3. Gas analysis techniques

The exhaust gas mixture was analysed via various analytical techniques, to ensure the accuracy of the measurements. Non-dispersive Infrared Spectroscopy (NDIR) (X-STREAM Enhanced General Purpose Process Gas Analyzer (XEGP) from Emerson) was used to analyse in real

$$Q_{CO_2out} (mol.min^{-1}) = 2 * Q_{CO_2in} / (2 * ([CO_2]_{out} / 100\% + [O_2]_{out} / 100\%) + [CO]_{out} / 100\%) \quad (5)$$

time the CO₂ and CO concentration profiles. Furthermore, we used a specific optical oxygen gas sensor (FDO₂, Pyroscience science technology) to measure the O₂ concentration continuously over time. We always started the measurements before plasma ignition, to first measure the blank composition, needed for the conversion calculation (i.e., 100% CO₂). All measurements were performed in triplicate.

In addition, we also analysed the exhaust gas with a three-channels compact gas chromatograph from Interscience. Each channel leads to a different column and detector. Two channels were used for this study: (1) carboxen and molsieve columns (1010 PLOT and 5A, respectively) were used to separate the permanent gases, such as O₂, N₂ and CO, further detected with a Thermal Conductivity Detector (TCD); and (2) two RT Q-bond columns (3 m and 10 m length, respectively) were used for the pre-separation and separation of CO₂ from the permanent gases, leading to a second TCD. The concentrations of the different products were calculated based on prior calibration of the GC in the concentration ranges of interest. As the GC analyses a sample immediately after sampling, a waiting time of about 6 min (time to get the chromatogram) was needed between one measurement and another. Moreover, this technique does not allow on-time measurements. Therefore, we compared the optical measurements (NDIR and optical O₂ sensor) with the GC measurements, to ensure their accuracy and to eventually use them for calculating the performance of the GAP with carbon bed in terms of CO₂

$$EC(kJ.L^{-1}) = SEI(kJ.L^{-1}) / (X_{corrCO_2} / 100\%) = (SEI(kJ.mol^{-1}) / V_m) / (X_{corrCO_2} / 100\%) \quad (8)$$

conversion and CO and O₂ concentrations.

2.4. Performance calculations

The conversion of CO₂ is usually defined as:

$$X_{CO_2} (\%) = ([CO_2]_{blank} - [CO_2]_{plasma}) / [CO_2]_{blank} * 100\% \quad (3)$$

With $[CO_2]_{blank}$ and $[CO_2]_{plasma}$ the CO₂ concentrations in the blank (without plasma) and during plasma, respectively (in %).

However, pure CO₂ splitting results in the expansion of the gas (CO₂ → CO + 1/2O₂) i.e., 1 molecule of CO₂ produces 1.5 molecules after (complete) reaction. In presence of a carbon bed, possible heterogeneous reactions between CO₂/O₂/O/CO and solid carbon might also occur, contributing even more to the gas expansion effect [34]. As a consequence, the exhaust gas flow rate is higher than the initial flow rate and all species are diluted in a higher volume than the initial one, i.e., the CO₂ concentration in the final volume appears lower than it would be in

the initial volume without expansion, resulting in an apparent higher conversion. Therefore, we need to know the exhaust flow rate, to correct the conversion from gas expansion. For this purpose, we used the method by Huang et al., [34]:

$$X_{corrCO_2} (\%) = (Q_{CO_2in} - Q_{out} * [CO_2]_{out} / 100\%) / Q_{CO_2in} * 100\% \quad (4)$$

With Q_{CO_2in} and Q_{CO_2out} the CO₂ flow rates at the inlet and outlet, respectively (mol.min⁻¹).

Q_{CO_2out} is determined based on the oxygen balance of the gas mixtures in the inlet and outlet [34]:

Besides the CO₂ conversion and the CO and O₂ concentrations in the exhaust mixture, we also calculated the energy efficiency and energy cost for our experiments. For this purpose, we need to know the specific energy input (SEI). The latter is the ratio of power over flow rate, and is defined as:

$$SEI(kJ.mol^{-1}) = (P * 60 / Q_{CO_2in}) * V_m \quad (6)$$

With P the power delivered into the plasma (kW), Q_{in} the inlet CO₂ flow rate (L.min⁻¹), 60 is the number of seconds per minute (s.min⁻¹) and V_m the molar volume at atmospheric pressure and room temperature (24.5 L.mol⁻¹).

The energy efficiency is defined as:

$$\eta (\%) = \Delta H_R^{\circ} * (X_{corrCO_2}) / SEI (kJ.mol^{-1}) \quad (7)$$

Based on our modelling insights, we expect the reverse Boudouard reaction to play only a minor role in the chemistry at the conditions under study. Therefore, the ΔH_R° used in the calculation refers to pure CO₂ splitting, i.e., $\Delta H_R^{\circ} = 283 kJ.mol^{-1}$ (cf. Equation (1)). See sections 4.1, 4.5 and 4.6 for further discussion.

Finally, the energy cost of the CO₂ conversion process is defined as:

3. Model description

We applied a zero-dimensional (0D) chemical kinetics model, using ZDPlasKin [45]. This model solves the continuity equations for the different plasma species.

$$\frac{dn_s}{dx} = \sum_j R_j [a_{sj}^R - a_{sj}^L] = \sum_j (k_j \prod_l n_l) [a_{sj}^R - a_{sj}^L] \quad (9)$$

Table 1
Species described in the model.

Gas-phase ground state species
CO ₂ , CO, O ₂ , O, C
Surface species
C(s), C(s)-O

^aC(s) and C(s)-O represent an active carbon site and an oxygen atom chemisorbed at an active carbon site, respectively.

in which n_s refers to the density of the species s , x is the position in the simulated reactor, index j refers to reaction j and index l refers to the different reactants of reaction j . α_{sj}^R and α_{sj}^L are the right- and left-hand side stoichiometric coefficients of species s , respectively, taking part in reaction j , k_j is the reaction rate coefficient, and $R_j = k_j \prod_l n_l$, is the reaction rate, with $\prod_l n_l$ being the product of densities n_l of species present on the left side of reaction j .

3.1. Chemistry set

The species in our model are listed in Table 1.

The concept of active surface site entails that heterogeneous reactions occur only at specific sites on which oxygen is adsorbed and forms complexes, which may ultimately lead to desorption of products, such as CO or CO₂. For carbonaceous materials, the active surface sites can be identified as structural features, such as defects in carbon layer planes and edge carbon atoms, disordered carbon atoms, heteroatoms (O,S,N) and mineral matter [46].

Prata et al. [47] applied the concept of active surface site to the modeling of the gas-surface reactions relevant for the ablative heat shield occurring during spacecraft atmospheric re-entry. We have chosen the chemical reaction set proposed by Prata amongst the other available sets [48–50], since it provides the most complete description for carbon oxidation reactions, along with a better agreement with experimental trends. The reaction set for the carbon oxidation is detailed in Table 2.

In this table, B is the active site density (in mole per square meter), k_B

Table 2
Heterogeneous reactions and their rate coefficients taken from Prata et al. [47] for solid carbon oxidation.

Reaction	Rate constant	Units
O(g) + C(s) → C(s)-O (O chemisorption)	$\frac{1}{4B} \sqrt{\frac{8k_B T}{\pi m_{O_2}}} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹]
C(s)-O → O(g) + C(s) (O desorption from mobile sites)	$\frac{2\pi m_{O_2} k_B^2 T^2}{N_A B h^3} e^{-44277/T}$	[s ⁻¹]
C(s)-O → O(g) + C(s) (O desorption from immobile sites)	$\frac{2\pi m_{O_2} k_B^2 T^2}{N_A B h^3} e^{-96500/T}$	[s ⁻¹]
C(s)-O + M → CO(g) + M M = O, O ₂ , CO, CO ₂ (Collision-induced CO desorption)	$\frac{1}{4B} \sqrt{\frac{8k_B T}{\pi m_M}} e^{-6000/T} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹] ^{a,b}
O ₂ (g) + C(s) → O(g) + C(s)-O (O ₂ partial chemisorption)	$\frac{1}{4B} \sqrt{\frac{8k_B T}{\pi m_{O_2}}} e^{-14500/T} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹] ^c
O ₂ (g) + 2C(s) → 2C(s)-O (O ₂ chemisorption)	$\frac{1}{4B^2} \sqrt{\frac{8k_B T}{\pi m_{O_2}}} e^{-8000/T} \frac{V}{A} \frac{10^{12}}{N_A^2}$	[cm ⁶ . s ⁻¹]
2C(s)-O → O ₂ (g) + 2C(s) (O ₂ desorption from mobile sites)	$\sqrt{\frac{N_A \pi k_B T}{2m_{O_2}}} 5 \times 10^{-5} e^{-15000/T} \frac{V}{A} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹]
2C(s)-O → O ₂ (g) + 2C(s) (O ₂ desorption from immobile sites)	$\sqrt{\frac{N_A \pi k_B T}{2m_{O_2}}} 1 \times 10^{-3} e^{-15000/T} \frac{V}{A} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹]
O(g) + C(s)-O → CO ₂ (g) + C(s) -O oxidation – only from mobile sites)	$\frac{1}{4B} \sqrt{\frac{8k_B T}{\pi m_{O_2}}} e^{-2000/T} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹] ^d
O ₂ (g) + C(s)-O → CO ₂ (g) + O(g) (CO ₂ desorption – only from mobile sites)	$\frac{1}{4B} \sqrt{\frac{8k_B T}{\pi m_{O_2}}} e^{-2000/T} \frac{10^6}{N_A}$	[cm ³ . s ⁻¹] ^d

^aactivation energy increased from 4000 K [47] to 6000 K, pre-exponential factor lowered from 100 [47] to 1 for the mobile sites and from 1000 to 1 for the immobile sites (see motivations in text).

^badded also CO and CO₂ as collision partners, according to the Eley-Rideal model described in Prata [47].

^cadded to the reaction set from the model by Zhlukov and Abe [49] (see text, not included in Prata [47], for studying its effects in the model predictions.

^dactivation energy increased from 500 K (Prata) to 2000 K, as suggested by Zhlukov and Abe [49] (see text).

is the Boltzmann constant [J.K⁻¹], h is the Planck constant [J.s], m_O , m_{O_2} , m_{CO} and m_{CO_2} are the O, O₂, CO and CO₂ mass in kg, respectively. Note that V/A is the ratio between the carbon bed volume and the carbon surface area and it is calculated according to the experimental conditions (hence: in our model here adopted from the data of Prata [47]).

In our study we keep the assumption of two different types of oxygen chemisorption, namely mobile for weakly bonded oxygen (single-bond character C-O) and immobile for strongly bonded oxygen (double-bond character C = O). Zhlukov and Abe [49] already introduced this distinction between at least two different types of adsorption into a kinetic model. This is also justified by several studies demonstrating the heterogeneity of the active surface sites in carbon materials. Lussow et al. [51] determined the active site density on Graphon samples with O₂ chemisorption and found that, above 400 °C, the oxygen saturation on the surface sharply increases, suggesting the presence of at least two types of active sites. Later on, Ahmed et al. [52] developed a kinetic model to describe the O₂ gasification of pyrolytic carbon, at low pressure (100 Pa) and in a narrow temperature range (748 – 898 K). The proposed mechanism involved the distinction between two types of active surface sites, and unveiled the importance of such distinction for an accurate description of CO₂ desorption. More recently, density functional theory (DFT) calculations demonstrated that the description of the heterogeneity of the surface of carbons through at least two types of active sites is needed to explain the desorption of CO₂ [53,54].

As suggested by Prata [55], we defined the distribution between sites giving mobile and immobile chemisorption to be 30% and 70%, respectively. The activation energy for the CO desorption stimulated by collisions with gas molecules (or atoms) is increased to 6000 K in order to match the experimental trends with our 0D kinetic model. For the same reason, we have lowered the pre-exponential factors to unity. The stimulated CO desorption is modeled as an Eley-Rideal reaction, where the pre-exponential factors typically take values between 0 and 1. Moreover, Prata [47] determined the activation energy and pre-exponential factors by fitting with molecular beam experiments, under the assumptions that i) the desorption products are immediately removed from the surface and ii) the hyperthermal O and O₂ sent to the surface accommodate to the surface temperature. However, products may also interact with the surface and the O/O₂ beam might not be in equilibrium with the surface temperature, also mentioned by the author [47]. The choice of a larger activation energy is also supported by the value proposed by Yang and Yang [56] for the spontaneous CO desorption (see Table 3), by the values estimated by Tremblay et al. [57] and by the values computed by Montoya et al. [58]. Following the same reasoning, we increased the activation energy also for the CO₂ desorption, by choosing a value proposed by Zhlukov and Abe [49] in their model for carbon oxidation. The experimental evidences from Panerai et al. [41] corroborate our choice. In particular, the authors show that CO₂ is produced only at temperature higher than 700 K under O₂ atmosphere, whereas the use of the energy barrier proposed by Prata would lead to overestimation of CO₂ for lower temperatures. Moreover, we performed a thermogravimetric analysis (TGA) (see Fig. S3 in SI) for charcoal 1, and the results are in good agreement with Panerai et al. i.e. CO₂ desorption onsets only above 700 K. Finally, it is worth to remark

Table 3
Heterogeneous Boudouard reactions and their rate coefficients taken from Yang & Yang [56].

Reaction	Rate constant	Units
CO(g) + C(s)-O → CO ₂ (g) + C(s) (Forward Boudouard)	$2.14 \times 10^9 e^{-16706.9/T} \frac{10^3}{N_A}$	[cm ³ .s ⁻¹]
CO ₂ (g) + C(s) → CO(g) + C(s)-O (Reverse Boudouard)	$9.07 \times 10^{10} e^{-24346.6/T} \frac{10^3}{N_A}$	[cm ³ .s ⁻¹]
C(s)-O → CO(g) (CO spontaneous desorption)	$\frac{2.14 \times 10^8 e^{-19323.6/T}}{R^* T}$	[s ⁻¹]

Where R is the ideal gas constant [L.atm.K⁻¹.mol⁻¹].

that the activation energy of the surface reactions varies with the structure of the carbon material, the oxidation degree of the surface and the presence of catalytic impurities.

Furthermore, to study the reactions occurring between CO₂ and the carbon active sites, we also implement the heterogeneous Boudouard reaction set proposed by Yang and Yang [56], reported in Table 3. The rate constants for the elementary steps were determined in a temperature range of 873–1173 K for graphite. In order to integrate it with the above carbon oxidation set, we assumed that the dissociation of CO₂ leads to formation of 30% mobile and 70% immobile oxygen chemisorption, as for atomic oxygen. The existence of at least two different types of oxygen adsorption is also documented for CO₂. For instance, Yang and Yang [56] found out that, after CO₂ gasification of graphite, the bond strength of the oxygen complexes formed on the monolayer edges is about 9 kcal.mol⁻¹ stronger than the complexes formed on the multilayer edges, and that they have an average bond character of two-thirds. Later, Calo and Perkins [51] stressed the need of accounting for the energetic heterogeneity of carbon surfaces in the kinetic analysis of the CO₂ gasification.

3.2. Experimental input parameters for the model validation

As mentioned in the Introduction, we have chosen to validate our model with the gasification experiments presented by Panerai *et al.* [41], separately under O₂ and CO₂ atmosphere. Indeed, their experimental setup has a simple design. It can be approximated as a plug flow reactor and allows to isolate the effect of O₂ and CO₂ in the oxidation and consumption of solid carbon. Moreover, the absence of a plasma further reduces the complexity of the process, and therefore the number of experimental uncertainties, and it will enable us to investigate in a second stage the effect of plasma in the gasification mechanism, establishing solid foundations for future model developments.

The experiments carried out by Panerai *et al.* [41] provide a set of experimental data which can be implemented into our 0D kinetic model with relative ease. A cylindrical glass tube is placed inside a furnace, within which the temperature can be assumed homogeneous and constant. The carbon bed is fitted inside the glass tube, approximately at the center of the furnace. The feed gas, either pure O₂, CO₂ or CO, flows through the tube and reaches the carbon bed heated by the furnace. The products are then analyzed downstream by Residual Gas Analysis (RGA) mass spectrometry. This simple experimental configuration can be approximated to a plug flow reactor under a limited number of assumptions, the most important being that the fluid is perfectly mixed in the direction perpendicular to the axis. In addition, the high-temperature oxidation experiments were performed using FiberForm, a carbon fiber material which is thoroughly characterized [59–61] and therefore an excellent candidate for modelling purposes. Thus, this experimental set of data represents a good benchmark for the validation of our 0D kinetic model. The parameters taken by the experimental work of Panerai [41] and used in our modelling study are listed in Table 4. Note that for the validation step with their experiments [41], the temperature is assumed constant and homogeneous throughout the reactor

Table 4
Parameters used for the 0D simulations, taken from Panerai *et al.* [41].

Test	Feed gas	T [K]	P [Pa]	C ₀ mass [g]	Carbon bed length [mm]	mass flow rate [g s ⁻¹]
O ₂ _A	O ₂	518	2172	1.366	20	2.36
O ₂ _E	O ₂	1086	3668	1.357	20	2.34
O ₂ _H	O ₂	1502	6180	1.372	20	2.34
O ₂ _H ₆₀₀	O ₂	1502	1900	0.372	7.18	2.34
CO ₂ _C	CO ₂	983	2674	1.350	20	2.31
CO ₂ _H	CO ₂	1413	4176	1.330	20	2.31
CO_C	CO	1508	3954	1.383	20	2.06

tube in the model. This assumption is needed since we had no information about the actual temperature at the carbon bed. We know from our experiments that the temperature increases locally where oxidation reactions take place, due to their exothermic nature. However, the use of a furnace, as in Panerai *et al.* [41], may help to maintain the temperature uniform in the carbon bed, and we believe that this supports our assumption. In our experiments, on the other hand, the carbon bed is partially heated by the plasma, and also locally heated by oxidation reactions. Therefore, we measured the temperature variations in the carbon bed, as well as without, with the aim of linking our modelling results, for fixed gas temperatures, to our experiments.

The notation for each individual test listed in Table 4 is taken from Panerai *et al.* [41]. The notation indicates the feed gas (O₂, CO₂ or CO), while the letters A, E, H and C were arbitrarily chosen by Panerai *et al.* to enumerate the tests, but in fact they correspond to different temperatures (see Table 4). Note that only the tests for which the product concentrations were recorded over time and provided from the experiments have been used in this study. For test O₂_H, we performed an additional simulation, named O₂_H₆₀₀, with pressure, carbon mass and carbon bed length after 10 min of treatment. For each single test, the temperature is assumed constant and homogenous throughout the reactor tube inside the furnace. The experiments were carried out at controlled mass flow rate, with variable pressure. The initial mass of each carbon bed sample was weighted, allowing for a more accurate calculation of the absolute total surface area (TSA), in square meters, for each individual test. We estimated a specific surface area of 0.16 m².g⁻¹ for FiberForm, following the calculations presented by Gusarov [62]. This estimation is in the same order of magnitude of other carbon fibers [63,64]. The TSA is a useful parameter to estimate the density of active sites C_{surf} (cm⁻³), as described in formula 10:

$$C_{surf} = b * B * N_A * TSA / V_{bed} * 10^{-6} \quad (10)$$

Where b is a fitting parameter, B is the active site density over the surface in mole m⁻², N_A is the Avogadro number and V_{bed} is the volume of the carbon bed in cubic meters. The choice of introducing the parameter b to calculate the active site density C_{surf} is justified by the influence of the experimental pressure on the reactivity of the carbon materials. Ismail [63] observed that at pressures above 2.67 kPa, a new type of active sites is activated. In molecular-beam studies, the pressure is typically near vacuum. Therefore, it is reasonable to assume that at higher pressures, such as in the experiments by Panerai *et al.*, the density of active sites is larger. This is consistent with the fact that without a fitting parameter to enhance the reactivity, the simulations predict a negligible effect of the carbon bed, in contrast with the experimental observations. We decided to keep b fixed at 30, as it provided the best agreement with the experiments. We do not expect b to vary significantly between one test and another, as the pressure is maintained within the same order of magnitude throughout the entire experimental study. We would like to point out that the chemistry set and the model developed for this study are only suitable to quantitatively simulate the performance of the gasification process under the conditions tested in the experiments by Panerai *et al.* [41]. With the aim of describing a different system (e.g. with different carbon materials or operational pressure), a tuning of the kinetic model may be required. Nevertheless, the model can still provide insights into the underlying mechanisms, even when the experimental conditions vary.

4. Results and discussion

4.1. Performance comparison: With and without carbon bed

Fig. 3 presents the maximum CO₂ conversion, and the corresponding energy efficiency and energy cost, without carbon bed (benchmark) and for two different carbon materials with different particle size, namely active charcoal 1 and 2 (see previous section). The obtained maximum

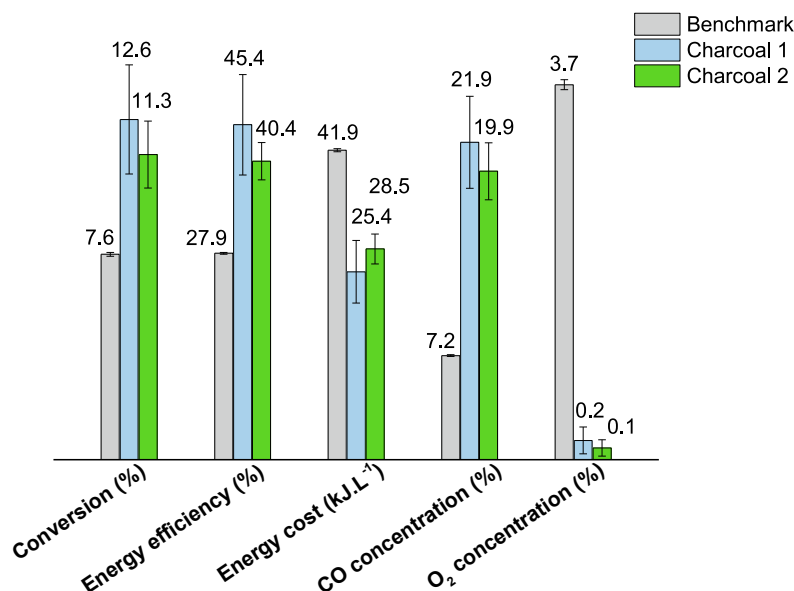


Fig. 3. Conversion, energy efficiency, energy cost, and CO and O₂ concentrations (obtained at the maximum conversion; see below), without carbon bed (benchmark), and with carbon bed (charcoal 1 and charcoal 2). 10 L_n.min⁻¹ CO₂. The SEI was about 3.2 kJ.L⁻¹ in all cases.

CO and O₂ concentrations are also depicted in Fig. 3.

It is known that the physicochemical properties of the carbon material and the particle size can seriously affect the carbon reactivity towards CO₂ during gasification [27,65]. More specifically, the concentration of active sites where the heterogeneous reactions between CO₂, O/O₂ (from gas) and carbon (from solid) occur is a critical parameter. Indeed, we observe a difference, with charcoal 1 producing slightly better results, but also somewhat larger error bars. This could be due to the larger pellet size, inducing more inhomogeneity in the bed. Nevertheless, we believe the better performance of charcoal 1 is attributed to a better gas and heat transfer. Indeed, larger particles leave more open space, which yields better convective heat transfer. As a result, when we measure the pressure upstream the carbon bed, the overpressure is < 550 Pa (=resolution of our measurements) with charcoal 1 and ca. 2800 Pa with charcoal 2. By contrast, charcoal 1, with larger pellet size, may limit the intra- and inter-particle diffusion, with a bigger effect on the former, affecting the reactivity of charcoal towards CO₂ [66]. However, because charcoal 1 and 2 have slightly different characteristics besides the particle size (e.g. charcoal 1 is steam activated, charcoal 2 is untreated), it is not possible to exactly estimate the effect of the mass transfer limitations. Indeed, charcoal 1, despite the larger particle size, provides higher CO₂ conversion and CO yield. We would like to address this question in our future work, where we will compare the performance of different pellet sizes, but of the same exact material, if available.

Nevertheless, whatever the type of charcoal (1 or 2), the conversion is significantly enhanced with the carbon bed, almost by a factor of two (from 7.6 ± 0.1 to 12.6 ± 2.0% and 11.3 ± 1.2% for charcoal 1 and 2, respectively), for the same SEI (about 3.2 kJ.L⁻¹). Consequently, the energy efficiency is increased from 27.9 ± 0.1% to 45.4 ± 6.8% and 40.4 ± 2.5% for charcoal 1 and 2, respectively, whereas the energy cost is much lower (from 41.9 ± 0.2 kJ.L⁻¹ for the benchmark to 25.4 ± 4.2 kJ.L⁻¹ and 28.5 ± 2.0 kJ.L⁻¹ for charcoal 1 and 2, respectively). The exact contribution of the RBR in the chemical process is not known, and, as mentioned in section 2.3, we expect it to play a minor role at the conditions under study (i.e., temperature in the carbon bed is only at the limit for the RBR to occur), so we calculated the energy efficiency based on ΔH_R⁰ for pure CO₂ splitting. This must however be kept in mind for the values obtained with the carbon bed. Finally, Fig. 3 also shows that the CO concentration is about three times higher upon addition of the

carbon bed, while the amount of O₂ is negligible when the carbon bed is present.

Therefore, adding a carbon bed not only favourably improves the CO₂ conversion and increases the CO concentration, but also gives much better energy efficiency and a much lower energy cost. In addition, the exhaust can be kept free of oxygen, which is very beneficial, as it avoids separation costs of CO from O₂. Indeed, besides removal and recycling of unconverted CO₂, CO purification is a technical challenge in case of scaled-up facilities, and thus an expensive process [37]. A CO₂/CO mixture obtained in the presence of a charcoal bed, free of O₂, could be directly recycled, by recirculating the gas in the plasma reactor, increasing the CO₂ conversion and eventually (ideally) leading to a pure CO stream. Note that recycling the exhaust in the plasma reactor again would reduce the overall energy efficiency, because energy is again needed for the second pass through the plasma. In real-life application, however, especially featuring near-zero cost renewable energy sources, there will be a trade-off between CAPEX (more reactors) and OPEX (recycling, hence more energy needed).

A similar improvement of the CO₂ conversion in the presence of a carbon bed was obtained with an atmospheric non-thermal plasma system studied by Huang *et al.* [34] (conversion up to 21.3% with carbon bed), but with a rather low associated energy efficiency (24.0%) due to a relatively high SEI (≈ 7 kJ.L⁻¹) compared with our conditions (where the SEI is 3.2 kJ.L⁻¹). In a thermal plasma, the CO₂ conversion was even more increased (up to 95 %) at a very high energy efficiency (70%), as studied by Li *et al.* [32]. However, these experiments were performed in CO₂/Ar mixtures, and Ar helps to enhance the CO₂ conversion and energy efficiency [8], but in reality, Ar also consumes part of the plasma energy, which was not accounted for when calculating both the conversion and energy efficiency. For a fair comparison, the fraction of CO₂ in the mixture should be accounted for [8,20]. Therefore, we cannot compare the absolute values. Moreover, because of the thermal plasma, the temperature was considerably higher (more than 1000 °C at 6 cm from the initial carbon surface) and therefore more beneficial for the RBR. A very high power (up to 16 kW) was necessary to drive this thermal plasma, which is far from our operating conditions (power around 0.5 kW). The authors proposed the following mechanism: CO₂ is first decomposed in a CO/O/O₂ mixture that reaches the carbon bed with a high temperature, favouring carbon oxidation by oxygen atoms and molecules; Second, the unconverted CO₂ is able to react with the

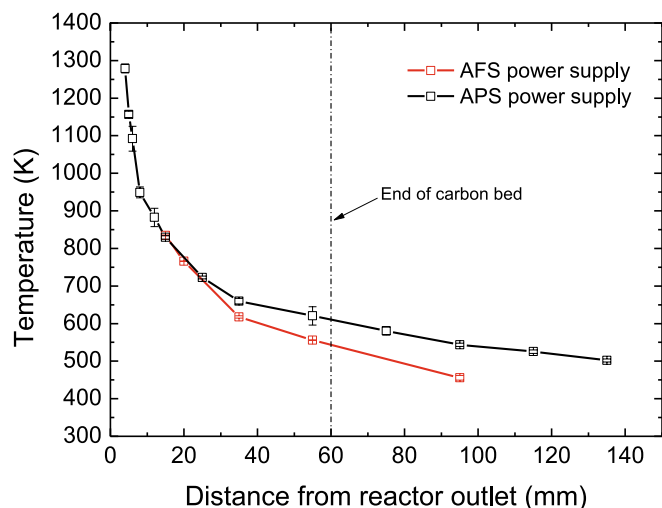


Fig. 4. Measured temperature profile in the post-plasma region without carbon bed (basket removed), with two different power supplies. Flow rate: $10 \text{ L}_n \cdot \text{min}^{-1} \text{ CO}_2$, SEI = 2.9 kJ.L^{-1} .

carbon bed via the RBR thanks to the elevated temperature from exothermic oxidation reactions [32].

In the next sections, we present the measured temperature profile after the plasma reactor, as well as the time-evolution of the CO_2 , CO and O_2 concentrations, to explain the influence of the carbon bed on the CO_2 chemistry after the plasma in our conditions. Furthermore, we will also present modelling results, to validate our hypotheses.

4.2. Temperature profile after the plasma reactor, without carbon bed

We measured the temperature profile after the plasma reactor, without carbon bed, to gain more insight in the chemical reactivity of the carbon towards CO_2 , CO, O and O_2 [41] (also needed to interpret the modelling results presented in section 4.6 below) and to determine the most suitable position of the carbon bed in the post-plasma reactor tube. The results are presented in Fig. 4 (red curve). Note that the post-plasma reactor tube was designed so that the temperature could be measured over the entire length (i.e. up to 255 mm), if necessary. However, we were most interested in measuring the temperature close to the reactor outlet, so we only measured up to 140 mm, because at further distances, the temperature is not so elevated anymore (see Fig. 4).

Note that the error bars, based on three independent measurements, are too small to be visible. Moreover, when the carbon bed is present, it starts at 0 mm (right at the reactor outlet) and extends till 60 mm (see vertical dashed line). In absence of the carbon bed, the temperature profile exhibits an exponential decrease with increasing distance from the reactor outlet (red curve). At 15 mm distance, the temperature reaches about 850 K and it drops to about 450 K at 95 mm distance. Note that for positions < 15 mm, the thermocouple tip was molten, as it was most likely in direct contact with the arc extending out of the plasma reactor, which made the measurements not possible. Based on previous modelling results, the arc temperature in the GAP is around 3000 K [44], so we can expect the temperature to be much higher than 850 K closer to the reactor outlet. Earlier, we had performed similar measurements for the same setup but with another power supply (APS; Advanced Plasma Solutions). However, the latter power supply later got broken, so it could not be used for our plasma conversion experiments. Nevertheless, we plot these temperature measurements in the same figure (black curve). They show a similar profile, and because the plasma arc behaviour was a bit different due to a different operational mode (attributed to the power supply), measurements closer to the reactor outlet were possible in this case, without melting the thermocouple tip. We could measure up to 1300 K at 4 mm distance from the reactor outlet. Although the

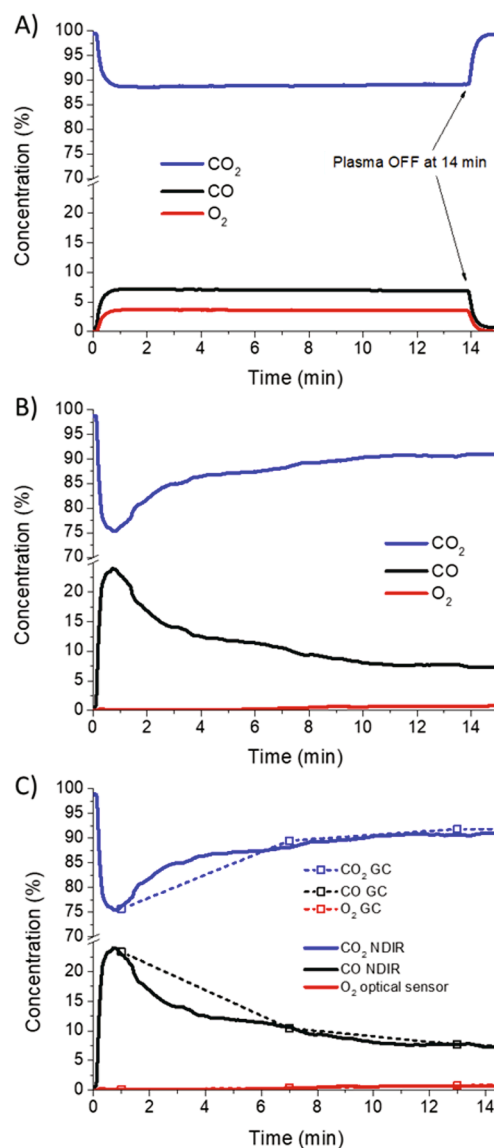


Fig. 5. CO_2 , CO and O_2 concentrations measured in real time, without carbon bed (A) and with carbon bed filled with charcoal 1 (B); C. Comparison of the optical sensors with GC measurements which only allow a few time points. Charcoal 1, $10 \text{ L}_n \cdot \text{min}^{-1} \text{ CO}_2$, SEI = 3.2 kJ.L^{-1} . The plasma is turned ON at time 0.

temperature profile is slightly different for a distance greater than 35 mm, the trend indicates a very similar behaviour for positions < 25 mm and therefore similar temperatures are expected close to the reactor outlet with the actual power source (red curve). Thus, positioning the carbon bed as close as possible to the reactor outlet (hence: to the plasma) allows the carbon bed not only to be in contact with the reactive O atoms from the plasma (providing quenching, so that they do not contribute to the back-reactions; see Introduction), but also to take benefit from the high temperature to enhance possible reactions between the exhaust gas and the solid carbon. This temperature profile will be used to explain the underlying chemical reactions, described in the model, in section 4.6 below.

4.3. Time-evolution of CO_2 , CO and O_2 concentrations, and associate CO_2 conversion

To better understand the chemical reactions occurring at the carbon bed and explain the very good results from Fig. 3, we plot in Fig. 5 the

concentrations of CO₂, CO and O₂ in the exhaust gas mixture as a function of time during plasma exposure, both without and with carbon bed (Fig. 5A and 5B) by using NDIR/O₂ optical sensors, and we also compare with GC measurements, to ensure a good accuracy of the optical sensors (Fig. 5C). Note, however, that the GC measurements only allowed a few time points; hence the advantage of the optical sensors. We only present the results for charcoal 1, but the data for charcoal 2 were very similar.

Without carbon bed (Fig. 5A), when the plasma is turned ON, the CO₂ concentration immediately drops to about 88%, whereas the CO and O₂ concentrations rise to 7.2% and 3.6%, respectively. In less than one minute, the three concentrations stabilize and stay constant over time, meaning that the plasma reaches an equilibrium, in line with previous studies with this plasma reactor [39,43]. Once the plasma is turned OFF, i.e., at 14 min, the chemical reactions stop and the concentrations return to their original values. The only reaction occurring here is CO₂ splitting, producing CO and O₂, as described in Equation (1) in the Introduction.

Upon addition of a carbon bed (Fig. 5B), the concentrations feature a quite different evolution over time. The CO₂ concentration immediately drops after plasma ignition and reaches a minimum value (ca. 75%) within 30–40 sec, and then gradually increases to ca. 90% at 14 min. The CO concentration reaches up to 24% when the CO₂ concentration is minimal, and then decreases again, similarly to the CO₂ increase. In contrast, the O₂ concentration depicts a very small transient rise right after plasma ignition ($\approx 0.2\%$, see Figure S8) before dropping to a negligible amount for about 5 min, followed by a very slow increase, but it stays below 0.8% as final value. A similar evolution was observed by Huang *et al.* with an atmospheric non-thermal plasmatron [34]. Therefore, the comparison between (A) and (B) shows that the carbon bed is deactivated after some time, for the production of CO (which relaxes to the benchmark value) but not for the O₂ consumption. The underlying mechanisms will be discussed later in sections 4.5 and 4.6.

To verify the accuracy of our NDIR and O₂ optical sensor measurements, we compared them with GC analysis, performed at the same time in parallel. GC is a very common technique for analysing gas mixtures and known to be very efficient and reliable; however, in the context of these experiments, on-time measurement is more useful to understand the complex chemistry. The results are presented in Fig. 5C and indicate a very good agreement between both techniques. Therefore, we can consider the on-time data reliable to calculate the conversion and the performance of the carbon bed.

From the time-profile of the CO₂ concentration (Fig. 5), we

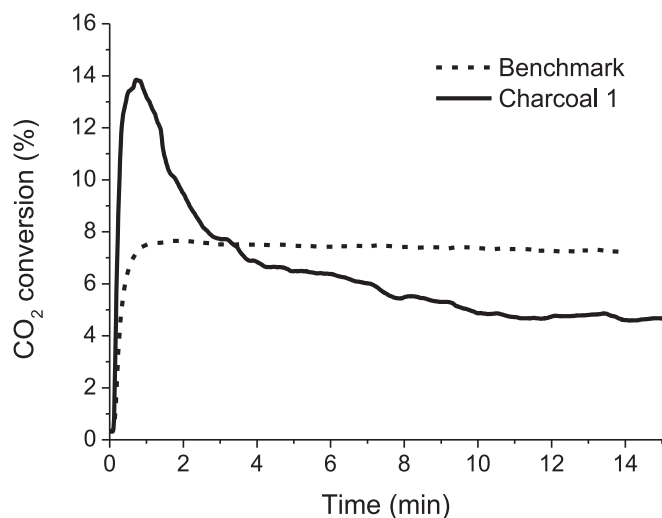


Fig. 6. CO₂ conversion as a function of time without (dash line) and with (solid line) carbon bed. Charcoal 1, 10 L_n.min⁻¹ CO₂, SEI = 3.2 kJ.L⁻¹.

calculated the CO₂ conversion (eq. (4) and (5)) for both the benchmark and the carbon bed addition, and the results are plotted in Fig. 6. In the first 1–2 min, the conversion upon carbon bed addition is much higher than without carbon bed, but it drops to lower values after ca. 3 min. In the next section, we discuss the results of the characterization study of the charcoals, before and after the gasification reaction.

4.4. Characterization of charcoal before and after the gasification reactions

The gasification of solid carbon is governed by interdependent factors such as the micro- and macrostructure of the material, the total and the active surface area, the degree of oxidation, the presence of catalytic impurities and the gas-phase composition. Any modification of one of these factors may alter the others. Therefore, the study of the underlying mechanisms for the gasification process, and their variations with time, cannot be separated from the screening of the modifications occurring during the reaction. For this purpose, we analysed the alterations of the surface exposed to the plasma-treated gas mixture, using SEM, N₂-BET and TGA-MS. The SEM images were collected for charcoal 1 as received and after 45 s and 7 min of reaction (see Figure S2 in SI). From the images, we could not identify any visible morphology change during the gasification reaction, except for the formation and deposition of ashes over the surface, which look more intense after 45 s rather than after 7 min of reaction. The BET measurements were performed on charcoal 1 and 2, as received, and after the gasification reaction for charcoal 1. The results are given in Table 5.

The BET measurements reveal that charcoal 1 and 2 have the same initial surface area. For this reason, we decided to perform the measurements after the reaction only for charcoal 1. After 45 s, the CO₂ conversion spikes (see Fig. 6) and the temperature is higher than 1300 K at 15 mm in the carbon bed, in the top layer, and still lower than 500 K at 35 mm, in the bottom layer (see Fig. 7, below). Therefore, it is reasonable to deduce that the gasification reactions are mainly occurring in the top layer of the bed. However, the surface area is barely affected, with a bigger effect on the top layer, most likely because of the desorption of volatile species (e.g. water) from the surface during the first seconds of reaction. After 7 min, the temperature in the top and bottom layer is nearly the same (see Fig. 7). At this point, Fig. 6 shows that the carbon bed is deactivated and the CO₂ conversion is lower than the benchmark value. Regardless of this, the surface area is further increased compared to 45 s. In view of the above, we cannot clearly link drop in performance of the carbon bed with a change in morphology or surface area. To investigate the presence of adsorbed species and thermolabile functional groups at the charcoal surface, before and after the gasification reaction, we carried out a TGA-MS study. The analysis highlighted that water is the main species desorbed by charcoal, already at low temperature (see Figure S4 in SI). This is an indication that water is mainly adsorbed over the surface and not coming from the decomposition of surface functional groups. Only at about 850 K, CO₂, CO and H₂ are also detected. We would like to point out that the analysis is not quantitative, and we need to be careful when comparing different products in terms of magnitude, because the mass spectrometer was not calibrated. However, at this point, H₂O desorption drops, and we hypothesize that a fraction of the other species detected are products of the reaction between water and charcoal, also called steam gasification [67,68]:

Table 5
BET measurements of the surface area.

Sample	BET surface area [m ² .g ⁻¹]
Charcoal 1	705 ± 71
Charcoal 2	703 ± 70
Charcoal 1, after 45 s, top layer	767 ± 77
Charcoal 1, after 45 s, bottom layer	732 ± 73
Charcoal, after 7 min	787 ± 79

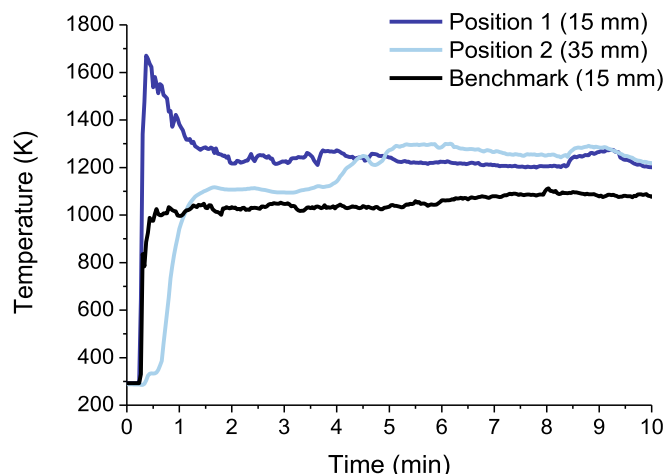
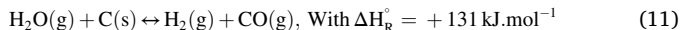
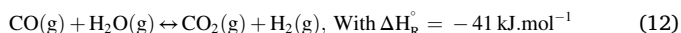


Fig. 7. Measured temperature at 15 mm and 35 mm from the reactor outlet, with carbon bed (dark and light blue lines, respectively), and without carbon bed (black line). Charcoal 1, $10 \text{ L}_m \cdot \text{min}^{-1} \text{ CO}_2$, $\text{SEI} = 3.2 \text{ kJ} \cdot \text{L}^{-1}$. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



The presence of CO_2 could be explained by either the water–gas shift reaction [68]:

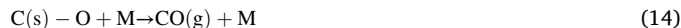
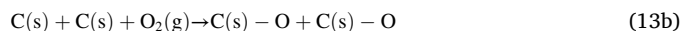


or by the decomposition of oxygen complexes or surface functional groups i.e. carboxyl groups [69], naturally found in charcoals. However, the TGA-MS of charcoal 1, after 45 s and 7 min of reaction (see Figure S5 and S6, respectively in SI), shows that water is still the main species desorbed from the surface, despite charcoal was certainly exposed to high temperatures during the gasification reactions. Such an observation indicates that water is probably re-adsorbed from the moisture in the atmosphere after the reaction, as charcoal is a hygroscopic material [70]. For this reason, it is not possible to draw clear conclusions from this analysis. However, the weight loss during the TGA (see Figure S7 in SI) shows that the charcoal tends to release less gaseous products with increasing gasification time, indicating that the surface may undergo a physico-chemical transformation that reduces the hygroscopicity of charcoal and/or modifies the nature or the amount of complexes and functional groups at the surface, favouring more stable interactions. With the results presented above, we hypothesize the underlying mechanisms in the next section.

4.5. Hypothesis of the underlying mechanisms

Based on the conversion and concentrations as a function of time, plotted in Figs. 5 and 6, we hypothesize that the following mechanisms take place:

1/ First, the conversion increases abruptly up to 14% (corresponding to the minimal CO_2 concentration observed in Fig. 5B). We can reasonably assume that, when the converted mixture (containing CO_2 , CO and O/O_2) is flowing through the carbon bed, O and/or O_2 are immediately adsorbed (either chemically or physically) on the carbon surface, as depicted in Fig. 5B and Figure S8 in SI. Hence, recombination reactions ($\text{CO} + \text{O}/\text{O}_2$ forming again CO_2) are suppressed and the net conversion increases, producing at the same time more CO compared with the benchmark (cf. Fig. 5A,B). The most likely mechanism occurring here is the partial oxidation of carbon by O/O_2 , following equations 13 and 14 [41,47,71]:



With $\text{C}(\text{s})$ a solid carbon active site. In Equations (13a,b), O and O_2 are immediately adsorbed by the carbon bed, followed by the release of a CO molecule promoted by collisions with gas molecules and high temperature (Equation (14)).

Note that for CO_2 splitting without carbon bed, the CO_2 conversion is usually equal to the CO concentration in the exhaust. Indeed, for the benchmark, both CO_2 conversion and CO concentration are about 7 – 7.5% in this case (Fig. 5A and Fig. 6). With carbon bed, the maximum conversion is 14% (Fig. 6), whereas the associated CO concentration is 1.5 times higher, i.e., 22%. Therefore, we can reasonably assume that part of the measured CO is released from the carbon surface, as described in Equation (14), in addition to CO coming from initial CO_2 splitting, which is most likely unreactive toward solid carbon [41]. A second hypothesis is that the RBR takes place as well (Equation (2) in the Introduction), but only if high enough temperature can be achieved (above ca. 973 K [27]) to activate the reaction towards the formation of CO. This will be further discussed in section 4.6 below.

2/ After this peak, the conversion decreases until reaching the benchmark value in around 3 min, whereas the CO concentration decreases and the O_2 concentration starts rising. Even more, from then on, the conversion continues dropping until being about 2% lower than the benchmark value, meaning that the positive effect previously observed does not last over time. We hypothesize that the carbon surface reaches saturation of oxygen: gas-phase O/O_2 is still being consumed by the remaining pellets that are not yet saturated (as seen in Fig. 5B and Figure S9) but it now leads to full oxidation and the release of CO_2 molecules (Equation 15). The role of the formation of oxygen complexes at the surface, most likely epoxy and peroxy complexes, on the CO_2 evolution during the carbon oxidation reactions, is also underlined by DFT calculations [72]. Moreover, the forward Boudouard reaction takes place (as demonstrated by our computer simulations; see section 4.6 below), explaining why the performance in terms of conversion becomes lower than the benchmark.



Moreover, as the oxidation reactions (sum of the elementary reactions 13 and 15) are exothermic [41,73,74], heat will be released and thus we expect an increase of the carbon bed temperature compared to the benchmark. In order to check this, we measured the temperature in the carbon bed (cf. Fig. 1 above) and compared it with the benchmark (in this case, the temperature is also measured in the basket, but free of carbon). The results are presented in Fig. 7, measured at 15 mm from the reactor outlet (dark blue line).

The temperature in the carbon bed reaches a maximum almost immediately after plasma ignition (ca. 1700 K), then decreases to finally stabilize around 1200 K after 2 min. Note that the temperature at 35 mm from the reactor outlet (light blue line in Fig. 7) is lower than at 15 mm during the first 4 min. This shows that the temperature profile is not uniform along the carbon bed and depends on the distance from the plasma. Moreover, we can assume that the carbon pellets at the end section of the basket are not yet saturated, as O atoms and O_2 are being consumed by the top layer. However, after 4 min the temperature at 35 mm increases and becomes higher than at 15 mm. It strengthens our hypothesis of full oxidation occurring at a later stage on the pellets placed further away from the anode exit (35 mm vs 15 mm) when they are saturated with oxygen (Equation 15).

As mentioned above, we believe that the initial rise in temperature at 15 mm, observed in Fig. 7, is due to exothermic reactions, most likely partial carbon oxidation (Equations 13–14) [75]. When the temperature is around 1200 K, CO_2 is most likely released from the surface, at the

expense of CO [41], and this is what we actually observe in Fig. 5B.

Moreover, it is very likely that the RBR occurs at such high temperature (1700 K) [32,34]. This is a highly endothermic reaction (see eq. (2) in the Introduction), so we can expect the temperature to be even higher than the measured one, but probably in a short period of time. Its contribution is however not straightforward to evaluate compared with carbon oxidation, which is more obvious from our experimental results. In any case, the presence of the carbon bed strongly influences the temperature, especially during the first minutes. Furthermore, the presence of the carbon bed might also affect the flow dynamics (physical effect), and consequently, this may affect the temperature. Hence, to discriminate between this physical effect and the chemical effects, we have replaced the carbon bed with an unreactive material (i.e., ceramic beads) that can withstand high temperatures. The temperature in the “dummy bed” did not exceed 1000 K, with associated CO₂ conversion and CO/O₂ concentrations similar to the benchmark values. Therefore, the main reason for the elevated temperature with the carbon bed is indeed most likely due to the exothermic reactions (13) and (15), as previously explained.

In comparison, the temperature in the basket without carbon (black line in Fig. 7) is lower and rather constant in the range of 1000–1050 K. This temperature is higher than at the same position without basket (about 850 K; see Fig. 4 above). This could be explained by the limited gas circulation inside the basket, its thick walls are heated up and they keep on releasing heat to the gas during the process.

In addition, the carbon bed may also represent a physical constraint to the flow, increasing the pressure upstream and causing a pressure drop throughout the bed. Panerai *et al.* [41] reported a pressure drop of up to ca. 90% under their conditions. Such a drop can be ascribed to the use of a roots pump to evacuate the system downstream their carbon bed. However, they did not discuss the effects of the pressure drop to the performance of their reactor. In our experiments, we could only estimate the drop by measuring the overpressure built upstream the carbon bed and assuming that the pressure downstream the bed is equal to atmospheric pressure. This is a good approximation since the exhaust flow is evacuated to open air, under the fume hood, without any pump. From our pressure measurements, we estimated the pressure drop to be lower than 0.4% for charcoal 1 and ca. 2% for charcoal 2. Therefore, we do not expect the pressure drop to play any significant role in the performance of the process under our operational conditions.

Finally, as shown in Fig. 1, the silo should provide fresh carbon once depleted in the basket and thus prolong the positive effect observed during the first minutes. Hence, the question arises whether the silo is not working properly. However, in another experiment in similar conditions (see Figure S9 in SI), we observed a sudden rise in temperature of about 100 K after ca. 10 min, followed again by a decrease. At the same time, the O₂ concentration drops to zero again, but for a very short period. This phenomenon is most likely due to fresh carbon supplied by the silo in the basket, leading to a drop in O₂ concentration and a small bump in CO₂ and CO concentration profiles (see Figure S9), following equations (9) and (10). However, the amount of fresh carbon seems insufficient to quench all the O/O₂, as observed in Figure S9 (the O₂ concentration rises again almost immediately, see inset figure), and therefore the positive effect remains limited. In any case, the temperature in the carbon bed seems very much linked to the CO₂, CO and O₂ behaviour in the mixture, but the cause and effect is not straightforward to understand and requires more in-depth analysis based on chemical kinetics modelling. Therefore, in the next section, we provide results of our model for the carbon gasification with different feed gases and we link the outcome to our experimental observations. This helps us indeed to better understand our experimental results.

4.6. Testing the hypothesis by means of modelling

Aiming to validate the set of reactions included in our model, we simulated the experimental tests listed in Table 4 in section 3.2 above,

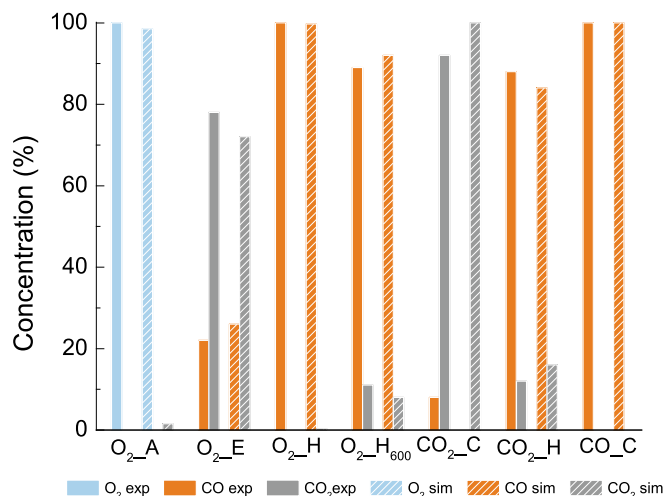


Fig. 8. Comparison between experimental (solid) [41] and simulated (diagonal pattern) product concentrations (see section 3.2 and Table 4 for the notation and the experimental conditions). The uncertainty on the concentration of the oxidation products was reported to be $\pm 12\%$ [41].

with the corresponding input parameters. The simulated product compositions, compared against the experimental values of Panerai *et al.* [41], are depicted in Fig. 8. Note that the product concentrations are taken at steady-state, which is reached within 200 s of treatment at all the conditions tested, except for O₂H₆₀₀, for which the comparison refers to the product concentrations after 600 s of treatment.

Fig. 8 shows that the model is able to capture the experimental trends and predict the product composition reasonably well. In particular, both our simulations and the experiments highlight no reactivity of solid carbon towards O₂ at 518 K (O₂_A), i.e., the only “product” is unreacted feed gas (O₂). Furthermore, they also agree on the total consumption of O₂ at higher temperature i.e. 1086 K (O₂_E) and 1502 K (O₂_H), producing 20–25% CO and 75–80% CO₂ at 1086 K, and pure CO at 1502 K. The reaction analysis for O₂_E and O₂_H (see Figure S10 and S11 in SI, respectively) indicates that O₂ promptly chemisorbs at the carbon surface, forming oxygen complexes. However, CO desorption is kinetically limited at 1086 K, therefore the oxygen complexes accumulate at the surface and can only undergo total oxidation (low activation energy), releasing CO₂ and consuming all the O₂. As a result, CO₂ is the main oxidation product in this case (O₂_E).

At 1502 K, CO desorption is kinetically favoured and competes with C(s)-O oxidation ($O_2(g) + C(s)-O \rightarrow CO_2(g) + O(g)$). Interestingly, the

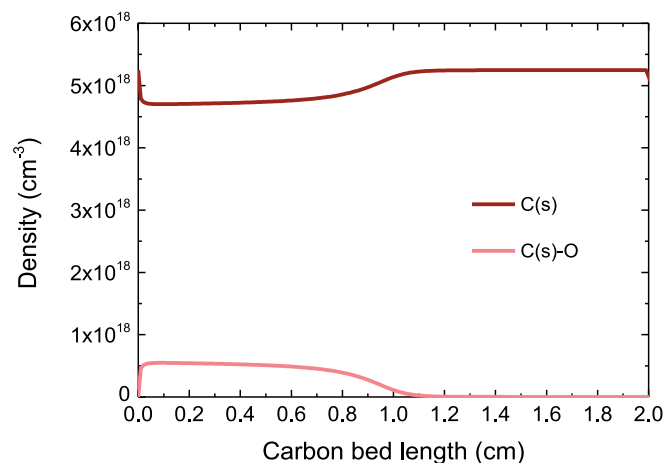


Fig. 9. Densities of free C(s) and occupied C(s)-O active sites as a function of the carbon bed length for the O₂_H test.

Boudouard equilibrium favours the CO₂ production until 1 cm in the carbon bed (see Figure S11 in SI). At that point, most of the active sites are free (see Fig. 9), as CO has been desorbed, and the RBR can proceed and convert CO₂ back to CO. As a consequence, O₂ is fully converted into CO, while CO₂ is not detectable as a product in this case (O₂H).

Therefore, we can divide the carbon bed in two zones: zone 1 where O₂ is depleted and carbon is primarily oxidized to CO₂, and zone 2 where CO₂ is consumed via the RBR. The length of the two zones depends on the temperature and the oxygen coverage (C(s)-O/C(s)) of the carbon surface. Indeed, despite the fact that the oxidation reactions typically occur within a few millimetres in the carbon bed (see Figure S10 and S11 in SI), the deactivation may start earlier than the substantial depletion of the pellets.

In order to get more insights into the deactivation mechanism, we simulated the gasification process after 10 min of treatment (O₂H₆₀₀). The experiments displayed that the carbon bed was no longer able to completely convert the CO₂ produced by the oxidation reactions through the RBR, due to both the reduced length of the bed (0.718 cm) and the saturation in oxygen of the surface. Indeed, in order to predict simultaneously a quantitative conversion of O₂ and the presence of CO₂ in the products, 11% of oxygen coverage at the surface has been added in the model. Without such addition, O₂ is not fully consumed by the oxidation of the oxygen complexes and only CO is produced. According to Fig. 9, we can reasonably assume that oxygen coverage builds up. The reaction analysis for O₂H₆₀₀ (see Figure S12 in SI) indicates that O₂ chemisorption rapidly slows down until O₂ is released instead, due to oxygen coverage at the surface. The onset of oxygen coverage with consequent release of O₂ favours the formation of CO₂. The Boudouard equilibrium is also shifted towards CO₂ production.

Thus, the deactivation of the carbon bed, as observed in our experiments after a few minutes (see Figs. 5 and 6 above) can be ascribed to the rapid formation of oxygen complexes at the surface, which undergo both fast oxidation to CO₂ and decomposition to CO. Both reactions efficiently consume solid carbon, reducing the bed size and the zone where CO₂ can react with the active sites. However, O₂ chemisorption is so fast that C(s)-O oxidation and decomposition are the limiting steps, leaving the surface partially saturated in oxygen. On top of that, the oxygen coverage hinders the sites for O₂ chemisorption, favouring oxidation, which promotes the forward over the reverse Boudouard reaction, producing additional CO₂ at the expense of CO. This explains the drop in CO₂ conversion and in CO concentration after a few minutes in our experiments (see Figs. 5 and 6 above).

Fig. 8 also shows that when pure CO₂ is introduced as input gas into the reactor tube, the performance of the carbon bed is constant over

time. In contrast to O₂, CO₂ is unreactive towards solid carbon at temperature lower than about 1000 K [27]. In fact, at 983 K (CO₂C) the production of CO is very small, resulting from the RBR. For similar conditions (O₂E, 1086 K), O₂ is promptly removed from the gas phase.

On the other hand, at 1413 K (CO₂H), the RBR slowly consumes CO₂, producing CO and oxygen complexes (see Figure S13 in SI). Such complexes release a second CO molecule upon collisions with gas molecules (or O atoms). Since the reaction, CO₂(g) + C(s) → CO(g) + C(s)-O, proceeds much slower than the O₂ dissociative chemisorption, the RBR is never hindered by oxygen coverage at the surface. Such a slower kinetics is also highlighted by the carbon consumption rates described in Panerai et al. [41]. As a result, at this condition more than 80% CO is produced, both in the experiments and predicted by our model (CO₂H).

Finally, CO was tested as feed gas and the experiments confirmed its inertness toward solid carbon, even at high temperature (CO₂C, 1508 K). Indeed, in absence of O/O₂ in the gas phase, oxygen complexes cannot be formed over the surface and CO cannot be oxidized through the forward Boudouard reaction. At the same time, CO does not present any reactivity towards the free active sites.

At this point, we validated our 0D kinetic model for the gasification of carbon under different atmospheres (O₂, CO₂, CO, separately) in a wide range of operating conditions. The insights achieved by this set of simulations, performed to reproduce the experimental results from Panerai et al. [41], help us to explain the underlying pathways for O/O₂ removal and CO formation, as well as for subsequent carbon bed consumption and deactivation occurring in our experiments. These mechanisms are schematically illustrated in Fig. 10, and can be summarized as follows.

When we turn on the plasma, a part of CO₂ undergoes splitting into CO and O. Subsequently, the O atoms will recombine into O₂ and/or react back with CO to form again CO₂. Therefore, the reactive mixture reaching the carbon pellets will be composed of a combination of CO₂, CO, O₂ and O. During the first millimetres of the carbon bed, O₂/O promptly chemisorb at the surface. The reaction is favoured both kinetically and thermodynamically, heating up the first layer of the pellets (see Fig. 7). The temperature spikes up to ca. 1700 K and CO desorption can proceed. At such high temperature, CO₂ can also be consumed through the RBR, increasing its conversion and further enhancing the CO production (see Figs. 5 and 6).

However, O₂/O chemisorption is faster than C(s)-O oxidation (O₂(g) + C(s)-O → CO₂(g) + O(g)) and CO desorption (C(s)-O → CO(g)) and the oxygen coverage at the surface starts to rise. Consequently, oxidation slows down, lowering the temperature, and O₂/O can reach the bulk of the carbon bed, explaining the increase in temperature further away

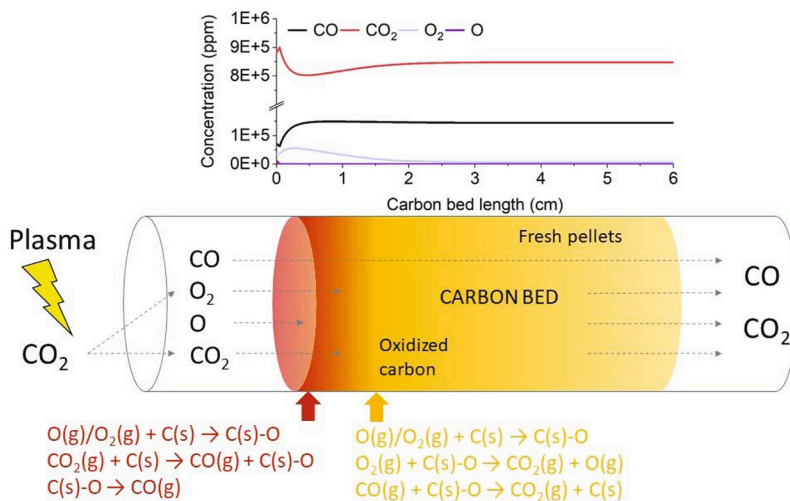


Fig. 10. Schematic illustration of the underlying mechanism for O/O₂ removal and CO (and CO₂) formation in the presence of a carbon bed after a CO₂ plasma.

from the plasma (see Fig. 7). The oxygen coverage deactivates the entire bed and CO₂ production is promoted by the full oxidation of the surface ($O_2(g) + C(s)-O \rightarrow CO_2(g) + O(g)$) and the forward Boudouard reaction ($CO(g) + C(s)-O \rightarrow CO_2(g) + C(s)$). At this point, the CO₂ concentration is higher than the benchmark, as well as the temperature, due to the exothermic oxidation reactions. Eventually, the oxygen coverage is so important that the carbon bed no longer quenches O₂, which can be found in the product composition. As mentioned, the overall mechanism for O/O₂ removal and CO (as well as CO₂) formation is summarized in Fig. 10.

Thus, the underlying mechanisms hypothesised based on our experimental observations are in line with the interpretation of our modelling results. In particular, the proposed mechanism demonstrates that the fast drop in performance of our experimental setup is not caused by a malfunctioning of the silo supplying fresh pellets. Instead, the operating temperature represents a crucial parameter, defining the composition of the gas-phase products and the rate of consumption of the pellets. Following these observations, we believe that providing additional heating to the carbon bed could help reducing the oxygen coverage and promoting the CO desorption. In this respect, an additional burner, most likely powered by gas, can purge periodically the amount of spent pellets left in the reactor, favouring the descent of fresh carbon from the silo. Another option can be the use of a thermal plasma, which can provide higher dissociation in the plasma region and, therefore, a more reactive mixture entering the carbon bed, with higher gas temperature, thus increasing conversion and energy efficiency [32,33]. However, as a drawback, argon had to be added to CO₂ in the feed mixture, in order to maintain the discharge and protect the cathode from carbon contamination [76]. The removal of argon from the product mixture would require additional energy thereby reducing the energy efficiency.

The other studies in literature mentioned earlier on the combination of plasma with carbon bed also reported that O/O₂ is rapidly and easily quenched by the carbon bed, although different carbon materials were used in the various papers, including ours. Note that we have a mesh between our plasma and carbon bed, which was not the case in [32,33]. Hence, the question arises whether this mesh could hinder the plasma coming in contact with the carbon bed. However, this is not the case, because in one plasma run the mesh was actually molten in its center, indicating that the plasma was even attracted towards the center of the mesh, as the latter is conductive. Moreover, if the mesh would cause an obstruction, leading to a lower peak temperature, we would observe that effect from the very beginning of the measurement, while this is not the case (see the experimental temperature profile in Fig. 7). Finally, in reference [34], Huang *et al* used an atmospheric non-thermal plasmatron with a stop-mesh on top of the carbon bed, to split the plasma into numerous micro-plasma jets after passing through the mesh. Therefore, adding a mesh at the anode exit is a way to attract the arc toward the outlet, directly towards the carbon bed. Huang *et al.* were able to observe a transient peak in the O₂ online-concentration profile 20 s after plasma ignition, that rapidly drops to zero [34], while Li *et al.* did not detect any oxygen in their experiments [32]. This is indeed very beneficial in terms of separation costs, as mentioned in section 4.1, in addition to the higher CO concentration obtained at these conditions. Hence, the key parameter to reach higher conversion appears to be the temperature, that depends not only on the type of plasma but also on the reactivity of the carbon bed toward O/O₂, and subsequently the ability of the carbon bed to release heat that can be directly reused to drive the RBR. Thus, future work should focus on searching the “ideal” combination of these parameters (type and size of carbon material, and operating conditions).

5. Conclusion

Plasma technology is very promising for CO₂ conversion, and especially warm plasmas, such as gliding arc and microwave plasmas, which exhibit good energy efficiency, but the conversion is subject to further

improvement. Moreover, CO₂ is converted into CO and O₂, which need to be separated in a second stage. In order to enhance the CO₂ conversion after the plasma, as well as the energy efficiency (thus reducing the energy cost), and at the same time promote O/O₂ removal and increase the CO fraction in the exhaust mixture, we placed a carbon bed after a Gliding Arc Plasmatron (GAP) reactor. Carbon pellets were constantly supplied by an innovative (silo) system, to avoid carbon depletion upon reaction with O/O₂.

This carbon bed allows to significantly enhance the CO₂ conversion, by almost a factor two (from 7.6% to 12.6%), while the energy efficiency rises from 27.9% to 45.4%, corresponding to a drop in energy cost from 41.9 kJ.L⁻¹ (without) to 25.4 kJ.L⁻¹ (with carbon bed). In addition, the CO concentration is about three times higher upon addition of the carbon bed, while the O₂ is nearly completely removed from the exhaust mixture, which is very beneficial, as it simplifies separation costs.

To understand the underlying mechanisms, we measured the temperature as a function of distance from the reactor outlet, and we also monitored the CO₂, CO and O₂ concentrations, as well as the temperature in the carbon bed, as a function of time. The time-resolved measurements reveal that the CO₂ conversion and CO concentration are only enhanced in the first minutes, followed by a drop to values below the benchmark (i.e., without carbon bed). To better understand this behaviour, we developed a model for the gasification of carbon under the effect of different gases (O₂, CO and CO₂ separately), which has been successfully validated with published experimental results.

Both our measurements and modelling results reveal that the onset of oxygen coverage at the carbon surface is of crucial importance to define the performance of the conversion process. In particular, the presence of oxygen complexes increases the selectivity towards CO₂ through the C(s)-O oxidation and the forward Boudouard reactions. In other words, the drop in CO₂ conversion and in CO concentration after a few minutes is attributed to the deactivation of the carbon bed, due to rapid formation of oxygen complexes at the surface.

This oxygen coverage can be limited by increasing the temperature in the carbon bed, i.e., by providing additional heating, leading to a complete consumption of the bed, followed by supplying fresh pellets from the silo. Hence, in our future work, we plan to apply additional heating, to further improve the performance. This additional heating will increase the energy cost of the process, linearly with the additional power supplied for the heating. This may lower the overall energy efficiency. However, the latter is not necessarily the case, if the rise in conversion due to the additional heating (promoting the specific chemical reactions) is higher than the increase in energy consumption. We aim to find conditions under which the beneficial effect of additional heating is larger than the additional energy cost. Nevertheless, our present results are already very promising, and clearly demonstrate the large potential of adding a carbon bed after a plasma reactor. In addition, our model provides very useful insights that explain our experimental observations and that are also useful for other experimental groups.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.cej.2022.136268>.

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